

# Proxy-SU(3)

Part I:

A focus on the highest weight irreps of the Elliott SU(3) symmetry

Part II:

Binding energies with the Proxy-SU(3) symmetry

SMARAGDA SARANTOPOULOU

INPP, DEMOKRITOS

This work has been funded by the NSRF (ΕΣΠΑ)

Smaragda Sarantopoulou, INPP, Demokritos  
18/5/21

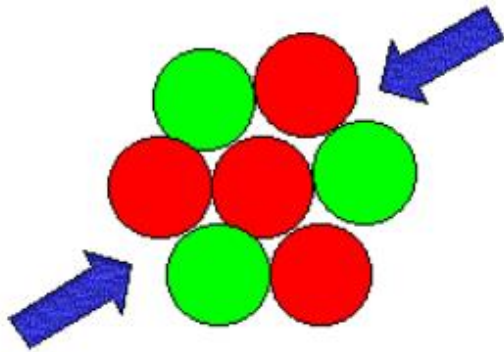


# Part I:

## A focus on the highest weight irreps of the Elliott $SU(3)$ symmetry

SMARAGDA SARANTOPOULOU  
INPP, DEMOKRITOS

# Basic properties of the strong force



**Strong force  
binds the nucleus**

- ❖ **Tensor force:**
  - Spin – isospin dependence
  - $d > 2 \text{ fm}$
  
- ❖ **Spin – isospin independent attraction:**  $\longrightarrow 1 \text{ fm} < d < 2 \text{ fm}$
  
- ❖ **Strong repulsive core:**
  - Pauli Principle
  - $d < 1 \text{ fm}$

# The many – particle wave functions

Antisymmetric  
many-particle  
wave function



Spatial wave  
function-most  
symmetric



Spinor-most  
antisymmetric

VOLUME 74, NUMBER 23

PHYSICAL REVIEW LETTERS

5 JUNE 1995

## Test of Wigner's Spin-Isospin Symmetry from Double Binding Energy Differences

P. Van Isacker,<sup>1</sup> D. D. Warner,<sup>2</sup> and D. S. Brenner<sup>3</sup>

<sup>1</sup>Grand Accélérateur National d'Ions Lourds, BP 5027, F-14021 Caen Cedex, France

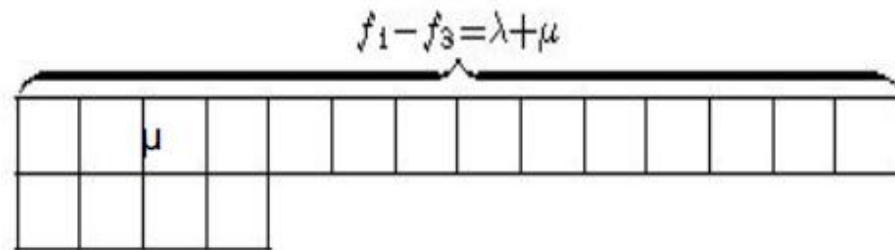
<sup>2</sup>DRAL Daresbury Laboratory, Daresbury, Warrington WA4 4AD, United Kingdom

<sup>3</sup>Clark University, Worcester, Massachusetts 01610

(Received 21 December 1994)

The most symmetric spatial wave  
function corresponds to the maximum  
Binding Energy.

# The Elliott SU(3) irreps



The Elliott SU(3) wave functions derive from  
The spatial part of the many-particle wave function

The Elliott SU(3) wave functions are many-quanta  
wave functions. Each box is a quantum in the z, x, or  
y cartesian axis.

The  $\lambda + \mu$  represent the quanta, which are symmetric  
upon their interchange

The  $\mu$  are not symmetric, neither antisymmetric



## An example

$$\begin{array}{|c|c|} \hline z & z \\ \hline x & \\ \hline \end{array}, \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 3 & \\ \hline \end{array}$$

$$\Phi_{spatial} = \sqrt{\frac{1}{6}} (2\phi_z(1)\phi_z(2)\phi_x(3) - \phi_z(1)\phi_x(2)\phi_z(3) - \phi_x(1)\phi_z(2)\phi_z(3)),$$

The quanta 1, 2 are symmetric upon their interchange, but 3 is not symmetric, neither antisymmetric

## Eur. Phys. J. A 57, 83 (2021)

$$r(\lambda, \mu) = \frac{\lambda + \mu}{\lambda + 2\mu} \cdot 100\%$$

The percentage of symmetric quanta over the total number of quanta gets the maximum value for the highest weight irrep of every shell.

valence particles	$\lambda$	$\mu$	$C_2$	$r$ (%)
1	3	0	18	100
2	6	0	54	100
3	7	1	81	89
4	8	2	114	83
5	10	1	144	92
	7	4	126	73
6	12	0	180	100
	9	3	153	80
	6	6	144	67
7	11	2	186	87
	8	5	168	72
8	10	4	198	78
9	10	4	198	78
	7	7	189	67
10	10	4	198	78
	7	7	189	67
	4	10	198	58

h.w.

h.w.

h.w.

h.w.

# Part II: Binding energies with the Proxy-SU(3) symmetry

SMARAGDA SARANTOPOULOU

INPP, DEMOKRITOS



# Method

## Elliott's Hamiltonian:

$$H = H_0 - \frac{\kappa}{2} Q Q$$

Where  $H_0$  is the 3D – HO Hamiltonian for the many - particle system:

$$H_0 = \sum_{i=1}^A \left( \frac{p_i^2}{2m} + \frac{1}{2} m \omega^2 r_i^2 \right)$$

If  $\hbar\omega = 1$  then:

$$H_0 = N_0 = \sum_{i=1}^A \left( \mathcal{N}_i + \frac{3}{2} \right)$$

With  $\mathcal{N}_i$  being the number of the harmonic oscillator quanta of the orbit of the  $i^{\text{th}}$  particle.

# Method

$$H = H_0 - \frac{\kappa}{2} QQ$$

Where:

$$\frac{\kappa}{2} = \frac{\hbar\omega}{4N_0} \quad \& \quad \hbar\omega = \frac{41}{A^{1/3}} \text{ MeV}$$

The quadrupole – quadrupole interaction in an SU(3) wave function is:

$$QQ = 4C_2 - 3L(L + 1)$$

$$C_2 = \lambda^2 + \mu^2 + \lambda\mu + 3(\lambda + \mu)$$

# Binding energy

$$BE(Z, N) = AV_0 \cdot \left( H_0 - \frac{\kappa}{2} QQ \right)$$

$$V_0 \approx 50 \left( 1 - \frac{N-Z}{A} \right) \text{ MeV}$$

Where  $V_0$  is the depth of the 3D – HO potential

The two proton or two neutron separation energies are given by:

$$S_{2p} = BE(Z, N) - BE(Z - 2, N)$$

$$S_{2n} = BE(Z, N) - BE(Z, N - 2)$$

# Calculation of $N_0$ in the HO shells

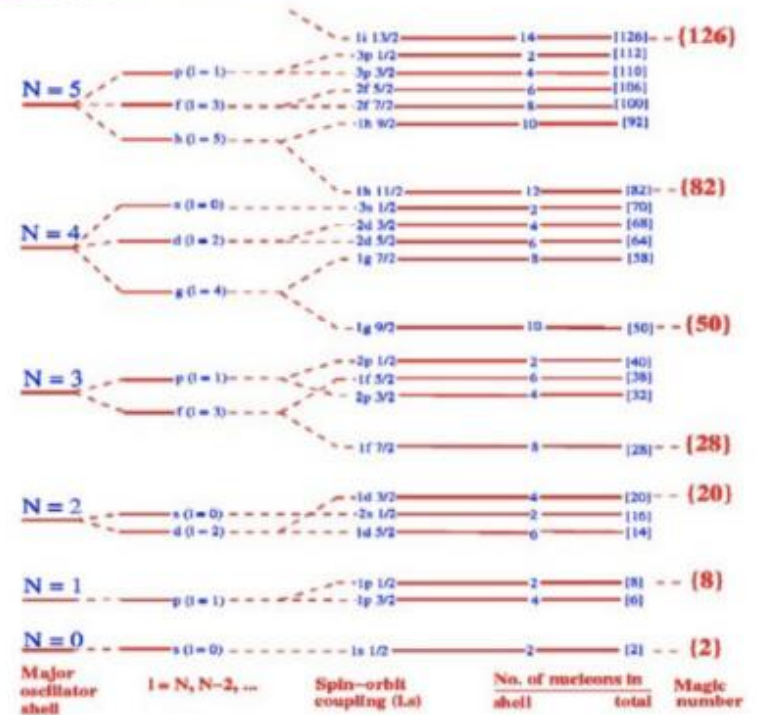


$$H_0 = N_0 = \sum_{i=1}^A \left( \mathcal{N}_i + \frac{3}{2} \right)$$

For the proton configuration:  $N_{0,Z} = 2 \left( 0 + \frac{3}{2} \right) + 4 \left( 1 + \frac{3}{2} \right) = 13$ .

For the neutron configuration:  $N_{0,N} = 2 \left( 0 + \frac{3}{2} \right) + 6 \left( 1 + \frac{3}{2} \right) = 18$ .

$$N_0 = N_{0,Z} + N_{0,N} = 13 + 18 = 31.$$



# Calculation of $N_0$ within the proxy-SU(3) symmetry

Within the proxy-SU(3) symmetry a unitary transformation can be applied in the intruder orbitals that reduces the number of quanta by 1. The proxy numbers become 6-12, 14-26, 28-48, 50-80, 82-124,...

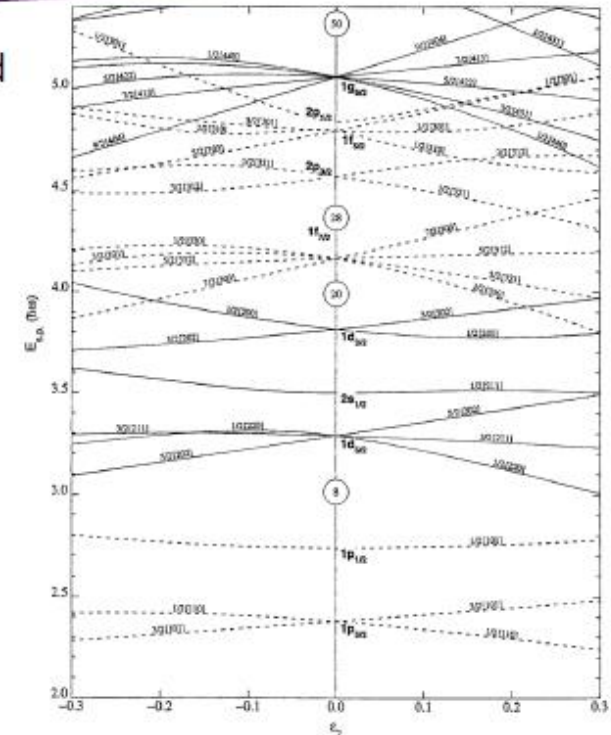
$$N_{0,SO} = N_{0,core} + N_{0,val}$$

$$Z=6 \quad N_{0,SO} = N_{0,core} = 13$$

$$N=8 \quad N_{0,SO} = 13 + (8 - 6) \left(1 + \frac{3}{2}\right) = 18$$

In general:

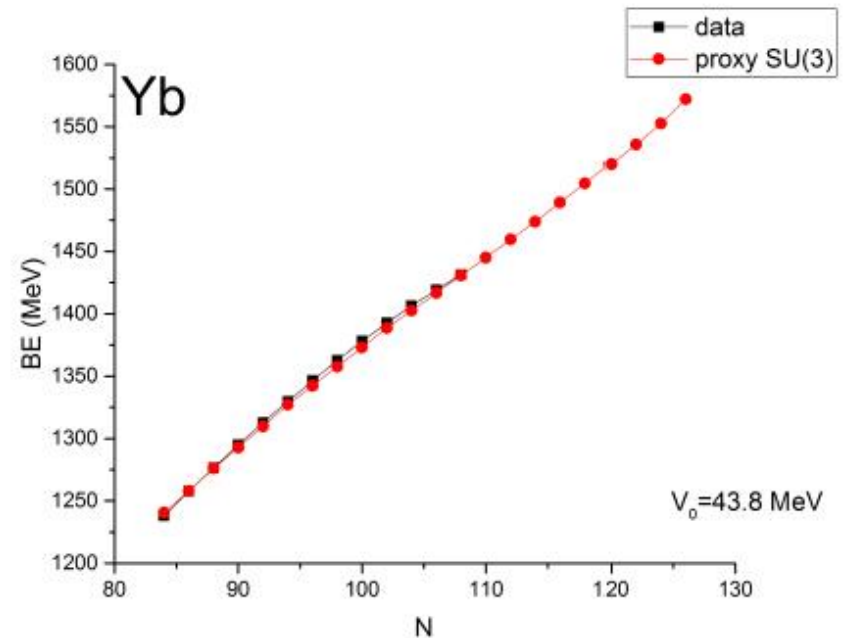
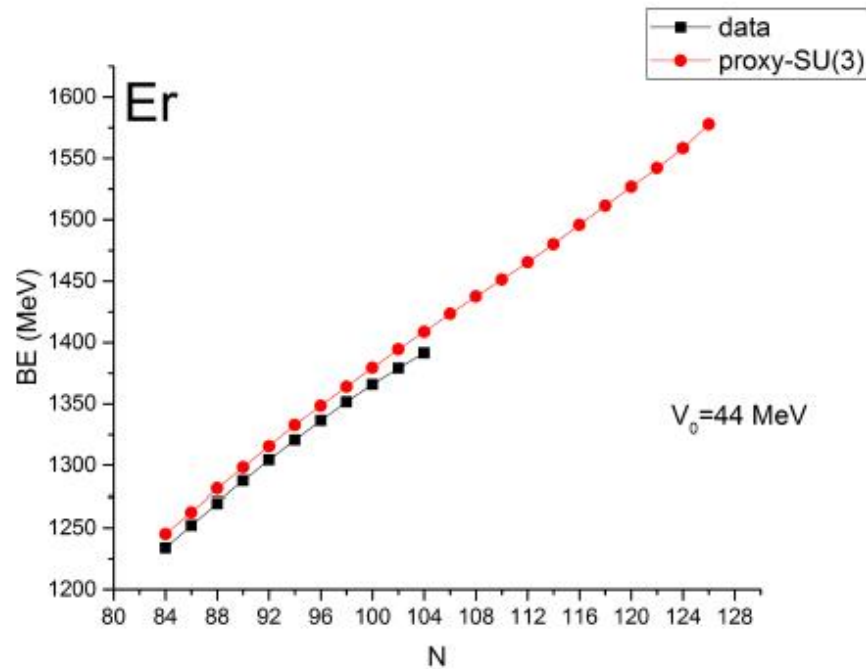
$$N_{0,SO} = N_{0,core} + \sum_{i=1}^{A_{val}} \left( \mathcal{N}_{i,proxy} + \frac{3}{2} \right)$$



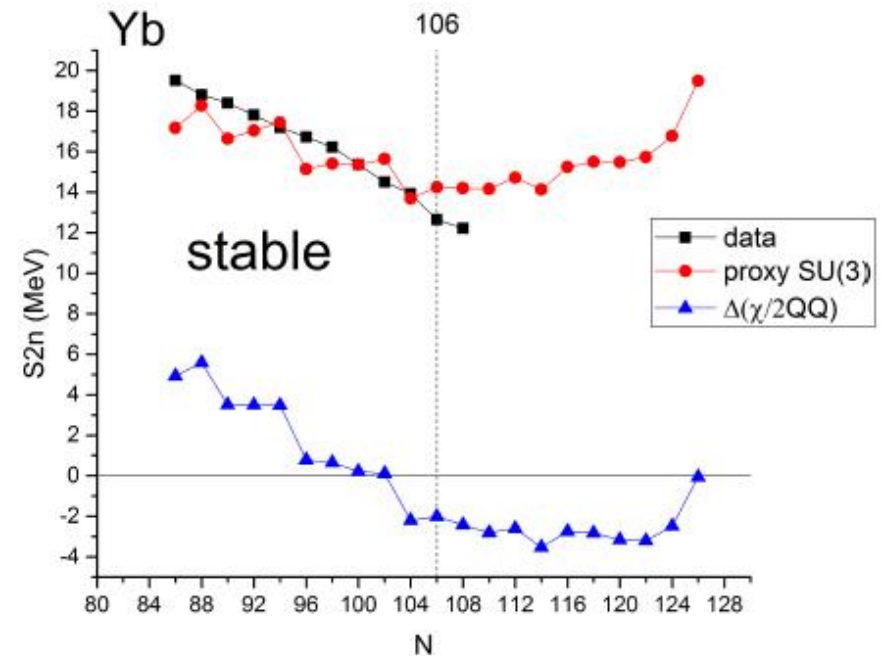
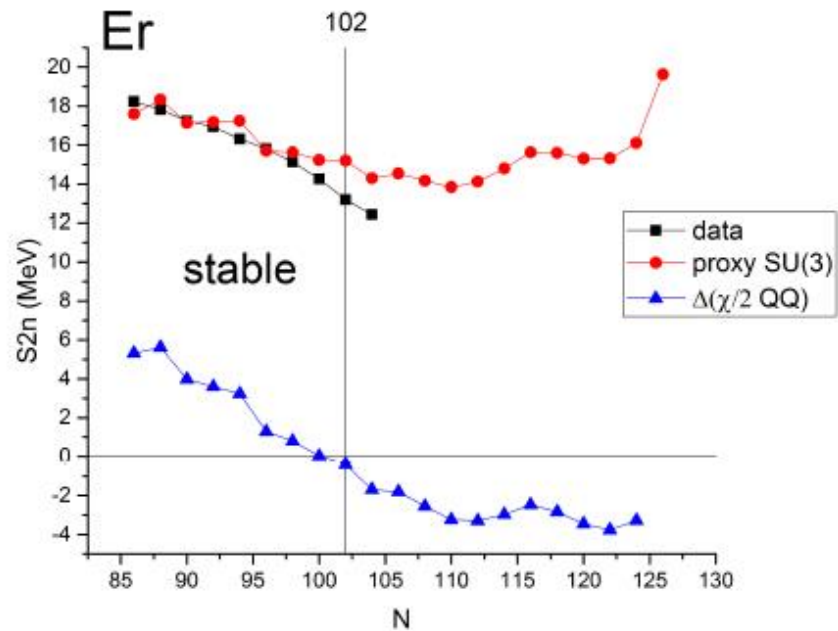


# Results with proxy-SU(3)

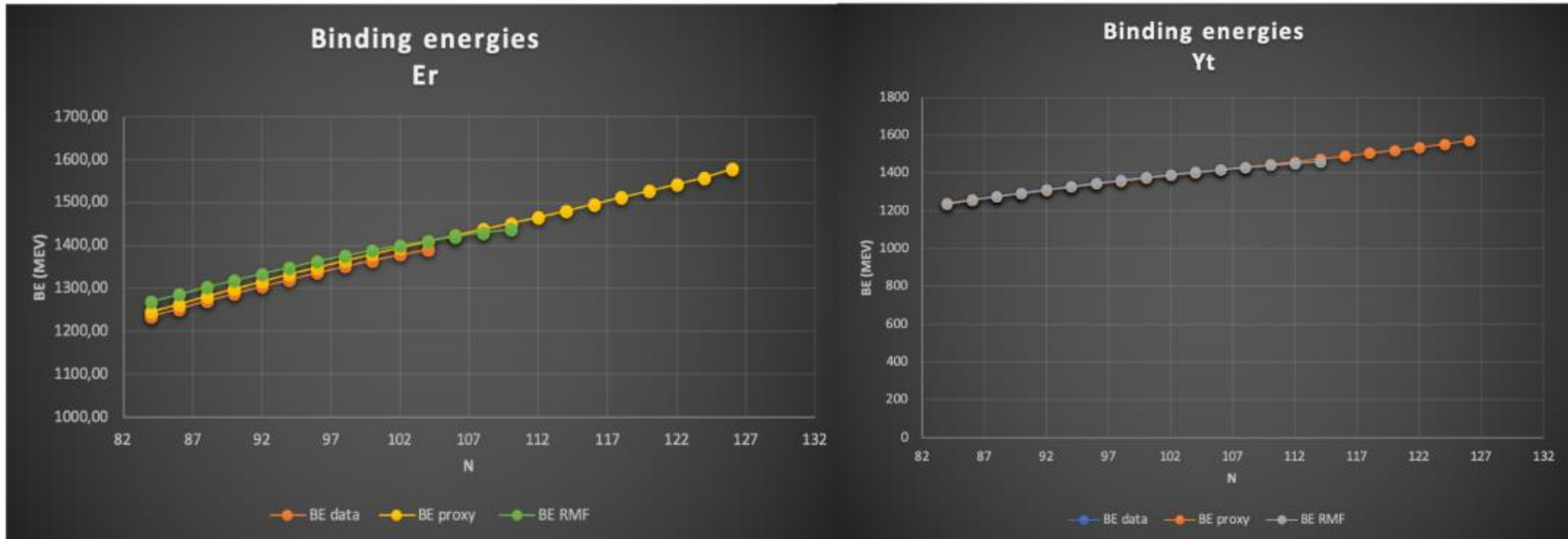
## Rare Earth nuclei



# Results with proxy-SU(3) Rare Earth nuclei



# Results in BE with proxy-SU(3) & RMF





# Extra



# The Elliott SU(3) irreps

${}^{96}_{42}\text{Mo}$

- $Z=42$  with 14 valence protons in 7 orbitals
- $[f]=[2,2,2,2,2,2,2,0,0,0]$

$$2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_{10}=14$$

$$2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_9=12$$

$$2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_8=10$$

$$2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_7=8$$

$$2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_6=6$$

$$2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_5=4$$

$$2 \ 0 \ 0 \ 0 \ \rightarrow S_4=2$$

$$0 \ 0 \ 0 \ \rightarrow S_3=0$$

$$0 \ 0 \ \rightarrow S_2=0$$

$$0 \ \rightarrow S_1=0$$

18/5/21

**The weight vector  $w$  is:**

$$W=(S_{10}-S_9, S_9-S_8, S_8-S_7, S_7-S_6, S_6-S_5, S_5-S_4, S_4-S_3, S_3-S_2, S_2-S_1, S_1)$$

$$W=(2,2,2,2,2,2,2,0,0,0)$$

# The Elliott SU(3) irreps

${}^{96}_{42}\text{Mo}$

- $Z=42$  with 14 valence protons in 7 orbitals
- $[f]=[2,2,2,2,2,2,2,0,0,0]$

**The weight vector  $w$  is:**

$$W=(S_{10}-S_9, S_9-S_8, S_8-S_7, S_7-S_6, S_6-S_5, S_5-S_4, S_4-S_3, S_3-S_2, S_2-S_1, S_1)$$

$$W=(2,2,2,2,2,2,2,0,0,0)$$

$$\begin{array}{r} 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_{10}=14 \\ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_9=12 \\ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_8=10 \\ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_7=8 \\ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_6=6 \\ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_5=4 \\ 2 \ 0 \ 0 \ 0 \ \rightarrow S_4=2 \\ 0 \ 0 \ 0 \ \rightarrow S_3=0 \\ 0 \ 0 \ \rightarrow S_2=0 \\ 0 \ \rightarrow S_1=0 \end{array}$$

# The many – body wave functions

$$H_i = \frac{p_i^2}{2M} + \frac{1}{2}M\omega^2 r_i^2 + V_{l_i s_i} \cdot l_i \cdot s_i$$

Nilsson orbitals

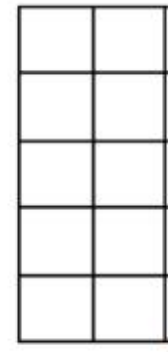
$K[Nn_z\Lambda]$	$(n_z, n_x, n_y)$
1/2[330]	(3,0,0)
3/2[321]	(2,1,0)
1/2[321]	(2,0,1)
5/2[312]	(1,2,0)
1/2[310]	(1,1,1)
3/2[312]	(1,0,2)
7/2[303]	(0,3,0)
3/2[301]	(0,2,1)
1/2[301]	(0,1,2)
5/2[303]	(0,0,3)



N=3

U(10)  
symmetry

Young diagram  
for 10 protons  
or neutrons



# The Elliott SU(3) irreps

${}^{96}_{42}\text{Mo}$

- $Z=42$  with 14 valence protons in 7 orbitals
- $[f]=[2,2,2,2,2,2,2,0,0,0]$

**The weight vector  $w$  is:**

$$W=(S_{10}-S_9, S_9-S_8, S_8-S_7, S_7-S_6, S_6-S_5, S_5-S_4, S_4-S_3, S_3-S_2, S_2-S_1, S_1)$$

$$W=(2,2,2,2,2,2,2,0,0,0)$$

$$\begin{array}{r} 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_{10}=14 \\ 2 \ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_9=12 \\ 2 \ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_8=10 \\ 2 \ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_7=8 \\ 2 \ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_6=6 \\ 2 \ 2 \ 0 \ 0 \ 0 \ \rightarrow S_5=4 \\ 2 \ 0 \ 0 \ 0 \ \rightarrow S_4=2 \\ 0 \ 0 \ 0 \ \rightarrow S_3=0 \\ 0 \ 0 \ \rightarrow S_2=0 \\ 0 \ \rightarrow S_1=0 \end{array}$$

# The Elliott SU(3) irreps

${}^{96}_{42}\text{Mo}$

**The weight vector  $w$  is:**

$$w = (S_{10} - S_9, S_9 - S_8, S_8 - S_7, S_7 - S_6, S_6 - S_5, S_5 - S_4, S_4 - S_3, S_3 - S_2, S_2 - S_1, S_1)$$

$$w = (2, 2, 2, 2, 2, 2, 2, 0, 0, 0)$$

For the weight vector the summations are:

$$f_1 = \sum_{i=1}^{10} n_{z,i} = 2(3 + 2 + 2 + 1 + 1 + 1) = 20$$

$$f_2 = \sum_{i=1}^{10} n_{x,i} = 2(0 + 1 + 0 + 2 + 1 + 0 + 3) = 14$$

$$f_3 = \sum_{i=1}^{10} n_{y,i} = 2(0 + 0 + 1 + 0 + 1 + 2 + 0) = 8$$

Young diagram:

$$[f] = (f_1, f_2, f_3) = (20, 14, 8).$$

The Elliott SU(3) irrep  $(\lambda, \mu)$  is given by:

$$\lambda = f_1 - f_2 \quad \& \quad \mu = f_2 - f_3$$

$$(\lambda, \mu) = (6, 6)$$

A  $(\lambda, \mu)$  irrep with  $\mu > 0$  corresponds to a wave function with mixed symmetry, while a  $(\lambda, 0)$  irrep to a totally symmetric spatial state.

We define the ratio:

$$r(\lambda, \mu) = \frac{\lambda + \mu}{\lambda + 2\mu} \cdot 100\%$$

which measures the percentage of the symmetric quanta  $\lambda + \mu$  out of the total number of quanta  $\lambda + 2\mu$ .