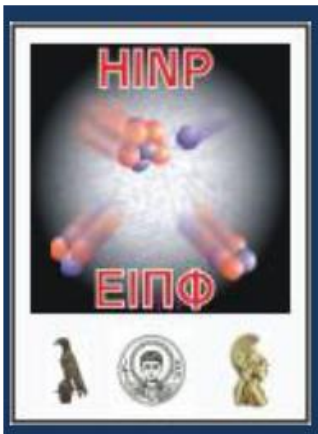




# Enhanced production of $^{99}\text{Mo}$ in inverse kinematics heavy ion reactions

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Prairie View A&M University



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6<sup>th</sup> Workshop of the Hellenic Institute of Nuclear Physics



## OUTILNE

- Introduction
- Experiment
- Data Analysis
- Results and Discussion
- Summary



# INTRODUCTION

**Methodology**



**Radioisotope production  
using inverse kinematics**

- Successful test of production of the theranostic radionuclide  $^{67}\text{Cu}$  ( $T_{1/2} = 62 \text{ h}$ )
- Reaction of a  $^{70}\text{Zn}$  beam at 15 MeV/A with a hydrogen gas target
- $^{67}\text{Cu}$  alongside other coproduced isotopes, collected after the gas target on an Al catcher foil and radioactivity measured by off-line  $\gamma$ -ray analysis.
- Use of secondary emitted particles from the primary nuclear reaction to irradiate other targets



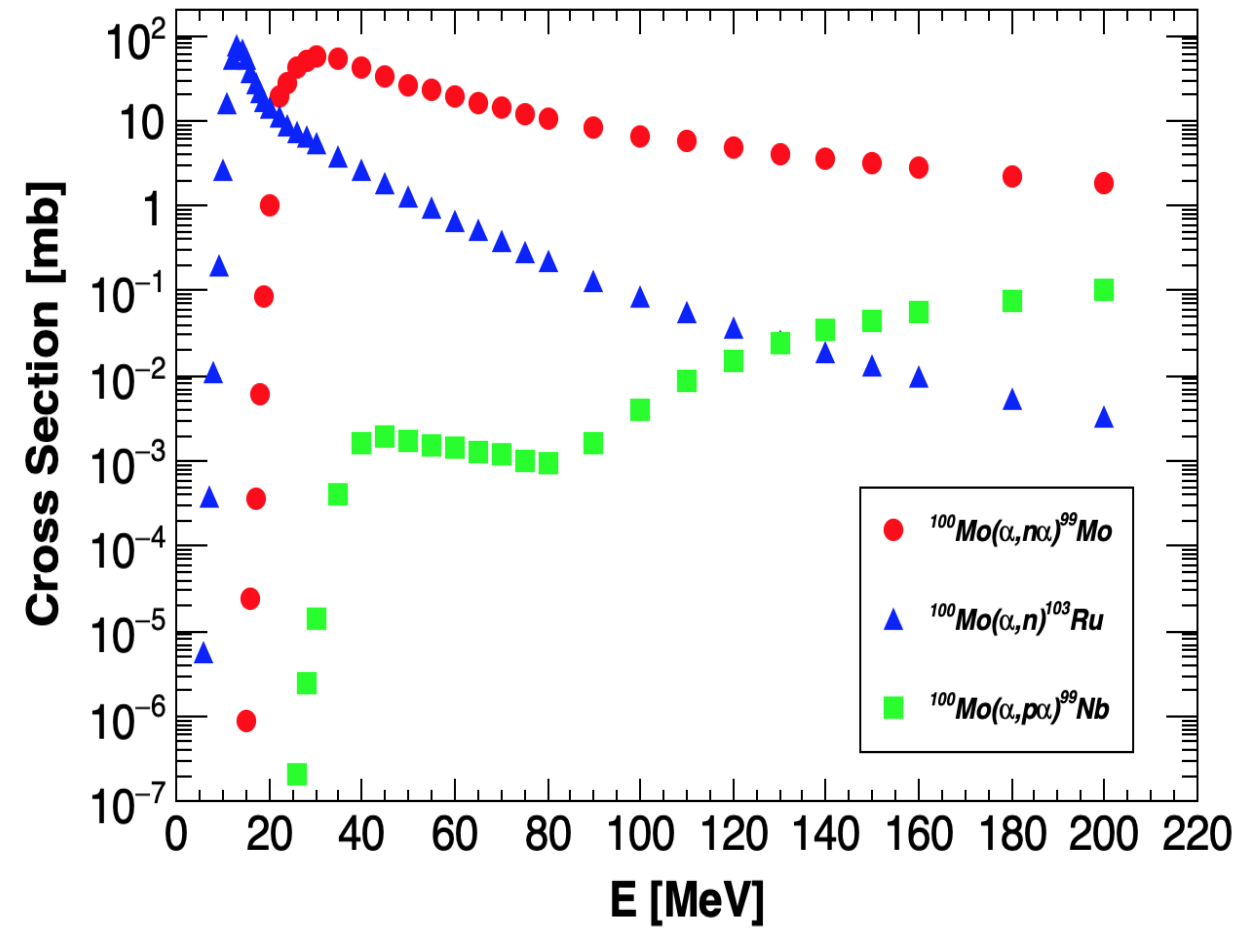
# INTRODUCTION

- Investigation of alternative production method for important isotopes for nuclear medicine in high demand worldwide
- Production in reactors not enough to supply the demand
- $^{99}\text{Mo}/^{99\text{m}}\text{Tc}$  generator: world demand estimated to  $\sim 7$  kCi/week
- $^{99\text{m}}\text{Tc}$ : 140 keV  $\gamma$ -ray emitter ( $I_{\gamma} = 89\%$ ,  $T_{1/2} = 6.01$  h) considered to be ideal radiotracer and used in  $\sim 85\%$  of all nuclear medicine diagnostic scans worldwide
- $^{99\text{m}}\text{Tc}$  is produced via  $\beta$ -decay from  $^{99}\text{Mo}$  ( $T_{1/2} = 65.94$  h)
- Other accelerator-made of  $^{99}\text{Mo}/^{99\text{m}}\text{Tc}$  :  $^{100}\text{Mo}(p,2n)^{99\text{m}}\text{Tc}$  and  $^{100}\text{Mo}(p,pn)^{99}\text{Mo}$



# INTRODUCTION

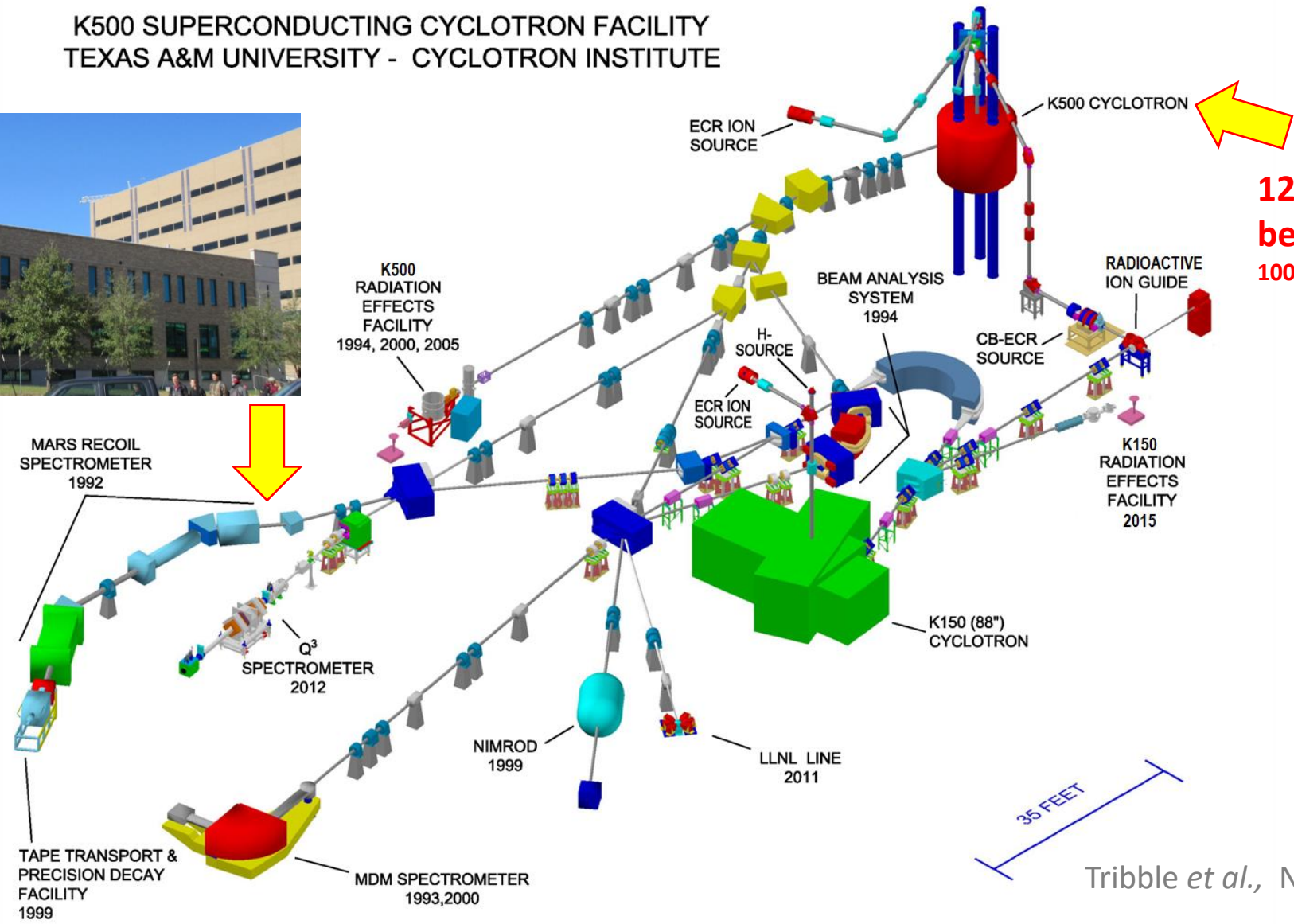
- Cross-section estimates from **TALYS** for different channels.
- Nucleon transfer on a  $^4\text{He}$  gas target produces resonances in  $^5\text{He}$  ( $^5\text{Li}$ ), and possibly high cross-section
- $^{100}\text{Mo} + ^4\text{He} \rightarrow ^5\text{He}$  ( $^5\text{Li}$ ) +  $^{99}\text{Mo}$  ( $^{99}\text{Nb}$ )  $Q = -9.2\text{MeV}$  ( $-13.1\text{MeV}$ )



[https://tendl.web.psi.ch/tendl\\_2019/tendl2019.html](https://tendl.web.psi.ch/tendl_2019/tendl2019.html)

# EXPERIMENT

K500 SUPERCONDUCTING CYCLOTRON FACILITY  
TEXAS A&M UNIVERSITY - CYCLOTRON INSTITUTE



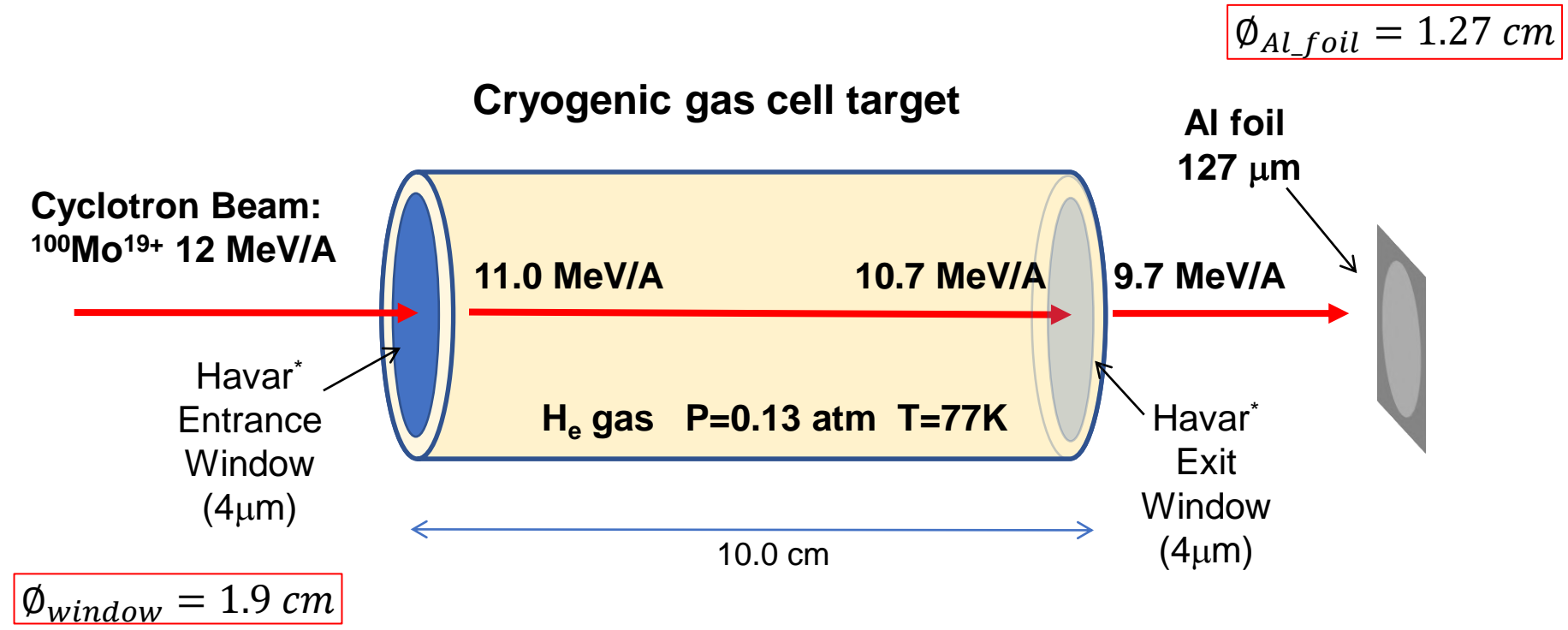
12 MeV/A  
beam of  
 $^{100}\text{Mo}_{19+}$

35 FEET

Tribble *et al.*, NIMA 285 (1989) 441-446



# Experimental Setup



\*HAVAR® - High strength non-magnetic alloy foil  
 Co 42.0% Mo 2.2% Cr 19.5% Mg 1.6% Ni 12.7% C 0.2% W 2.7% Fe Balance

Souliotis *et al.*, Appl. Radiat. Isot. 149 (2019) 89-95



# Experimental Specifications

## Irradiation

- $^{100}\text{Mo}$  (12 MeV/A) at an average charge state of  $35^+$
- 3 different runs
  - **Source 1** (102 torr, 0.07 pA, 11h35)
  - **Source 2** (213 torr, 0.21pA, 10h28);
  - **Source 3** (1008 torr, 0.17 pA, 7h52)
- Beam current was periodically monitored and was nearly constant (within 15 %)

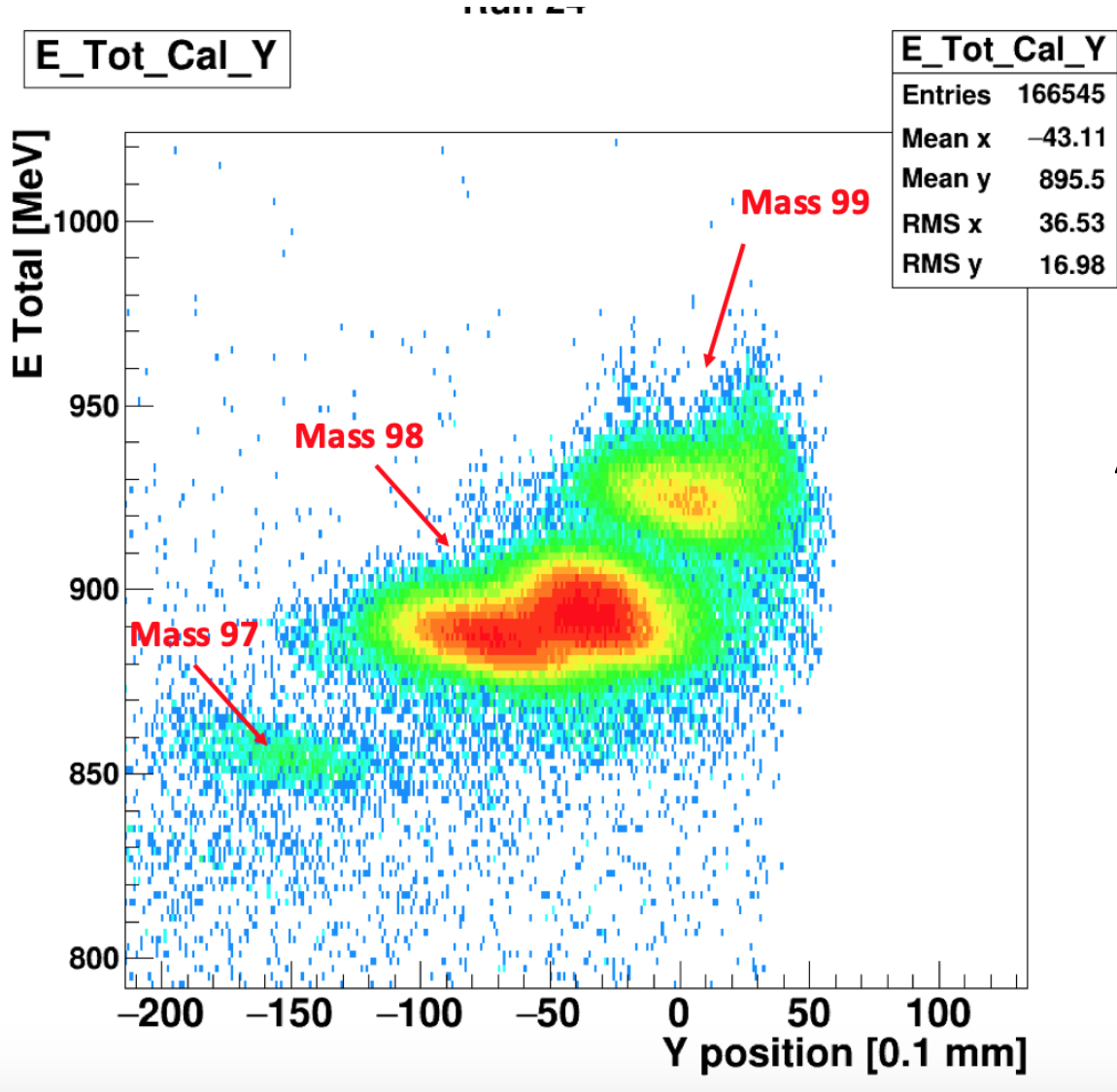
## Off-line $\gamma$ -ray analysis

- HPGe detector
- Al foil @  $d = 50$  mm
- dead time 2-3%
- Energy resolution 2.5–4.0 keV (FWHM)





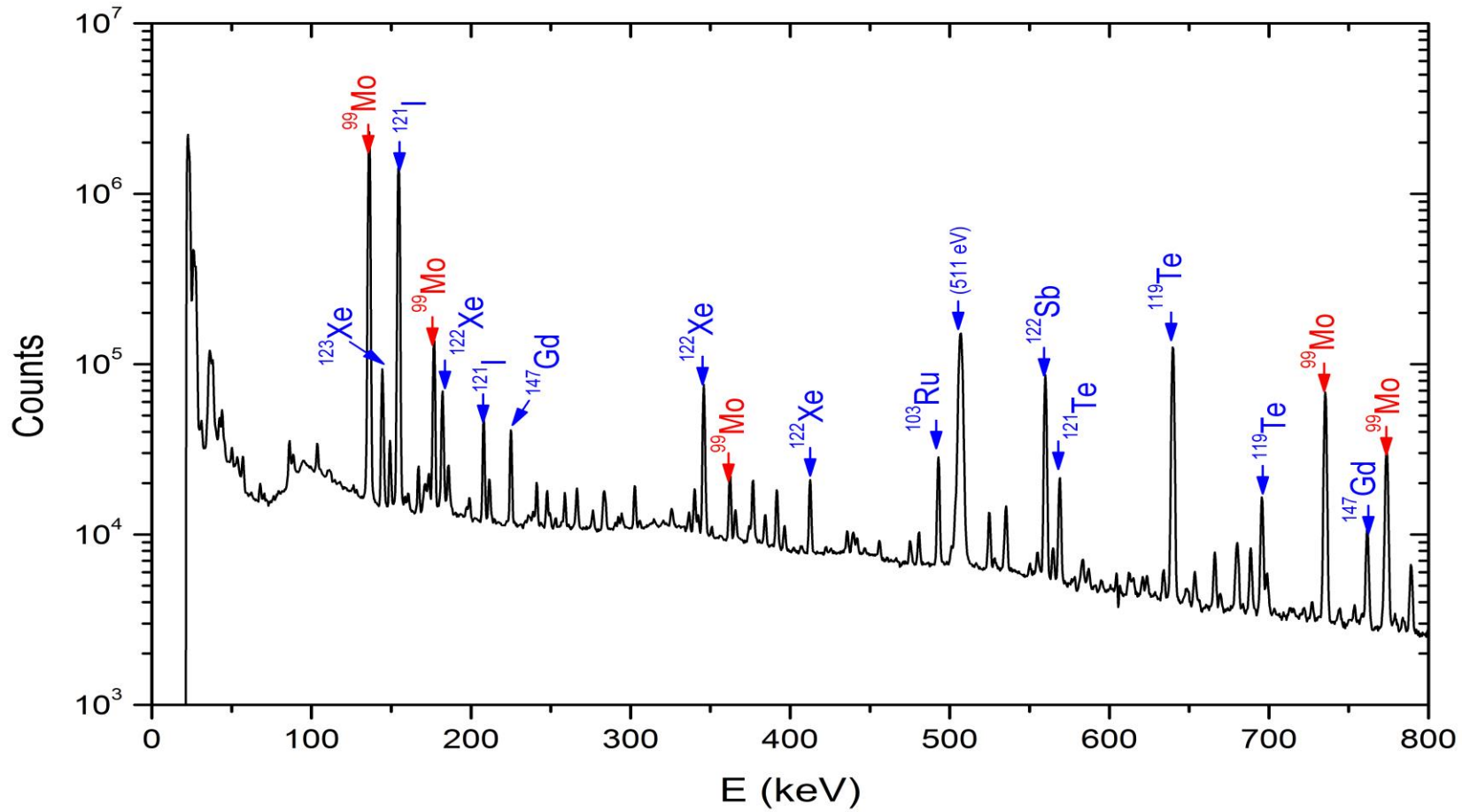
# DATA ANALYSIS



Focal plane of MARS  
 $\Delta E - E$ , Si detector 55  $\mu\text{m}$ -500  $\mu\text{m}$



# $\gamma$ – spectra from from source 3





# Activity after irradiation

Averaged Observed Activity

$$R_M = \frac{C}{\epsilon \Delta t_L I_\gamma}$$

$C$  → Net area of the peak  
 $\epsilon$  → Efficiency  
 $\Delta t_L$  → Measurement time  
 $I_\gamma$  → Branching Ratio

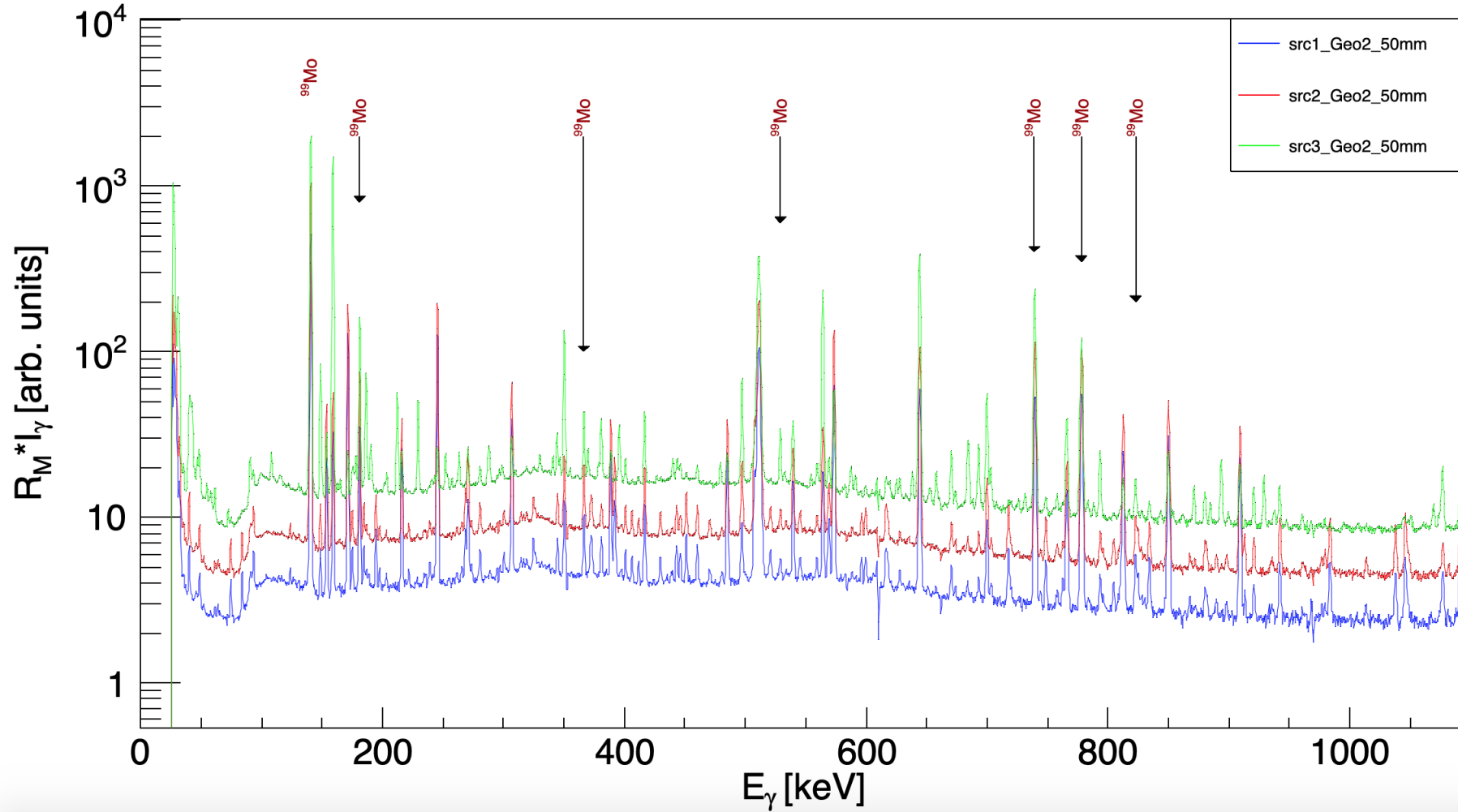
Activity after Irradiation

$$R_{AI} = R_M \underbrace{e^{\lambda t_{dk}}}_{\text{Correction for activity drop during Cool time}} \underbrace{\frac{\lambda \Delta t_R}{(1 - e^{-\lambda \Delta t_R})}}_{\text{Correction for decay during measurement}} \underbrace{\frac{\lambda t_{irr}}{(1 - e^{-\lambda t_{irr}})}}_{\text{Correction for decay during irradiation}}$$

$\lambda t_{dk}$  → Cool time  
 $\lambda \Delta t_R$  → Real time  
 $\lambda t_{irr}$  → Beam time  
 Decay Constant



# $\gamma$ – spectra from 3 runs with different He gas densities





## Activity after irradiation

$$R_{AI} = R_M e^{\lambda t_{dk}} \frac{\lambda \Delta t_R}{(1 - e^{-\lambda \Delta t_R})} \frac{\lambda t_{irr}}{(1 - e^{-\lambda t_{irr}})}$$

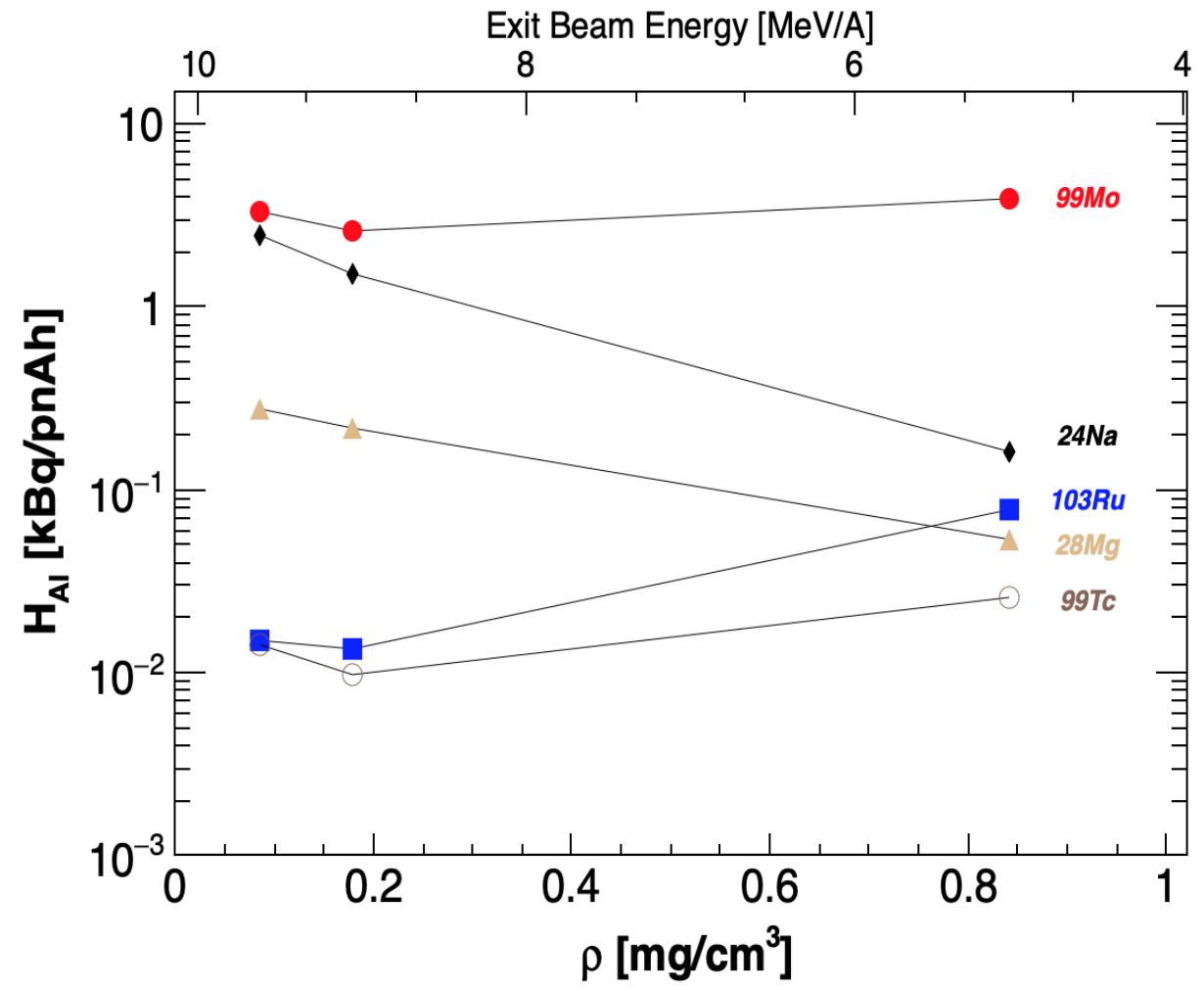
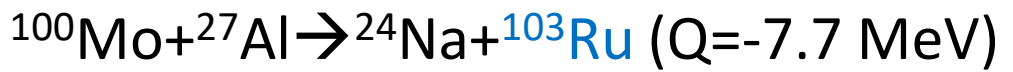
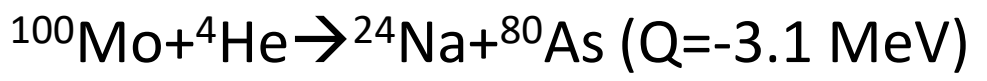
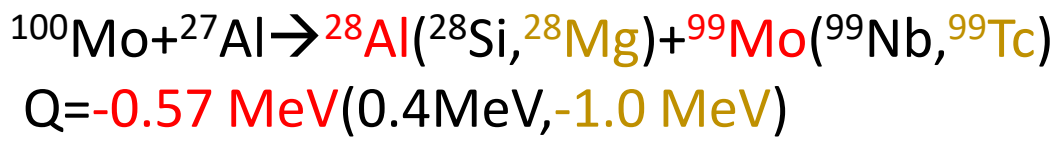
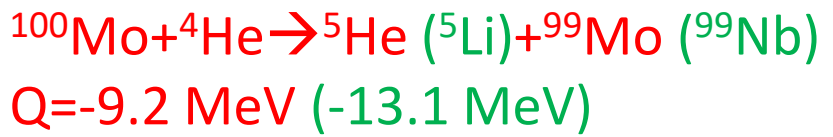
**Net Activity** = Activity produced per beam intensity unit and irradiation time

$$H_{AI} = \frac{R_{AI}}{I_{beam} t_{irr}}$$



# Activities for different isotopes of interest

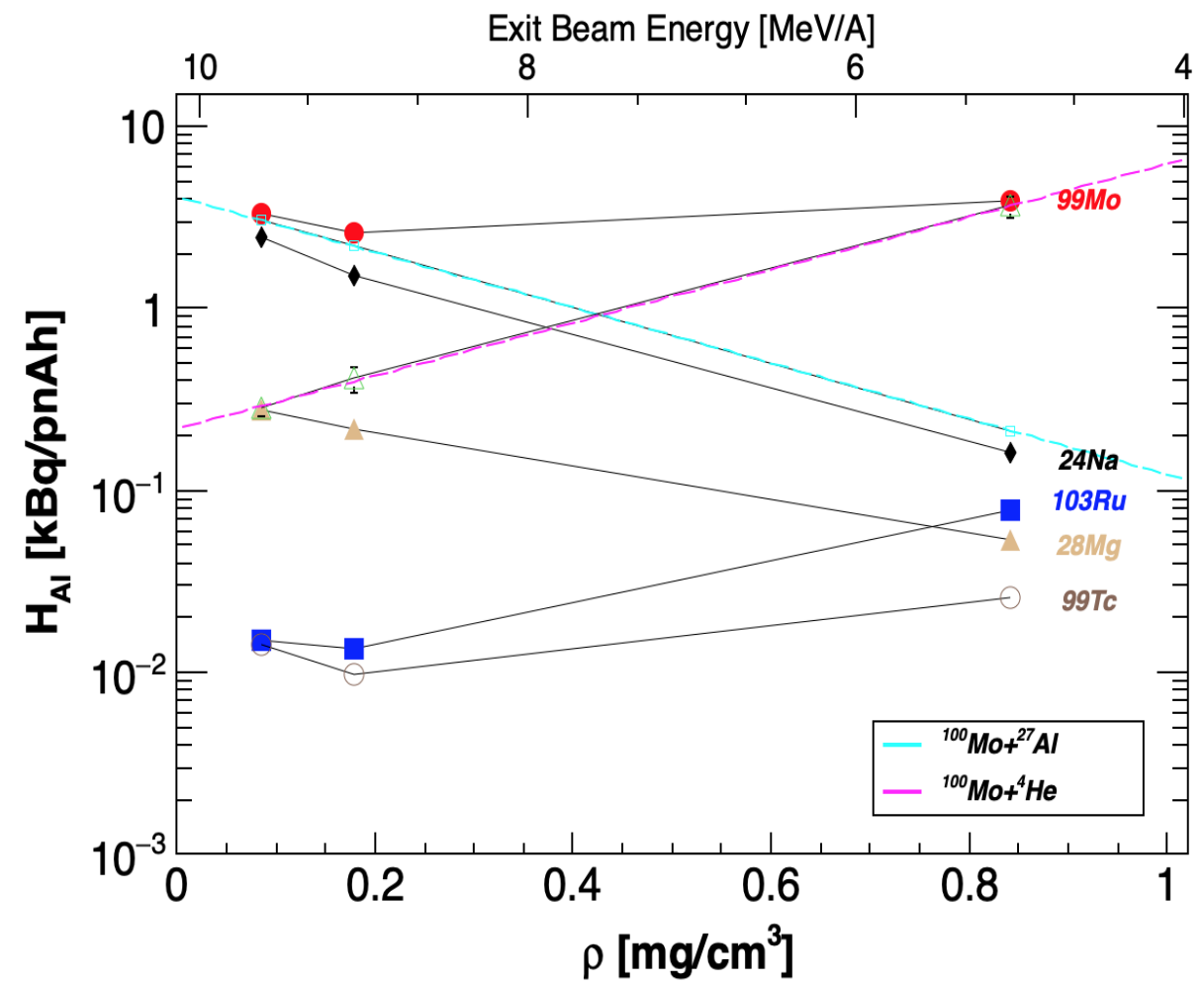
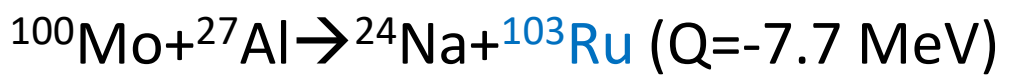
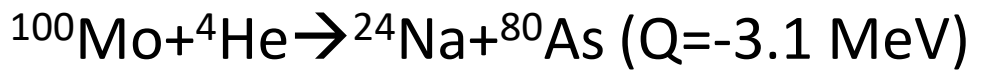
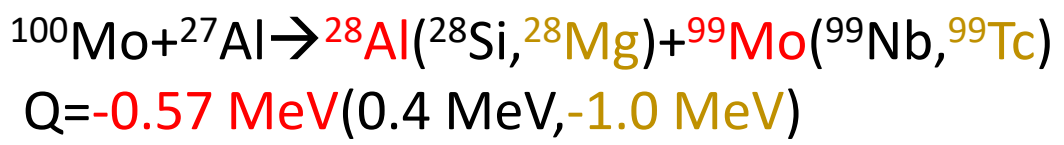
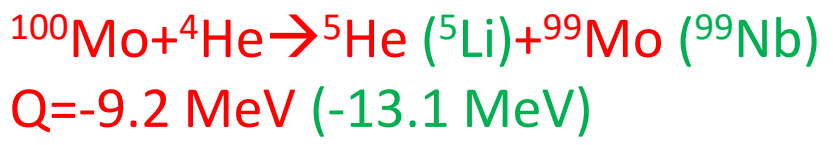
- Identified using their strongest and independent  $\gamma$ -lines





# Activities for different isotopes of interest

- Identified using their strongest and independent  $\gamma$ -lines





## SUMMARY

- This work has demonstrated the feasibility and advantages of using inverse kinematics to produce  $^{99}\text{Mo}$  and other isotopes.
- A gas target and a catcher have been used for production.
- Neutrons produced in the forward direction could be used to irradiate other targets and produce a variety of radioisotopes of interest.
- Production could also be determined for neutrons in coincidence with heavy fragments detected in MARS.
- Currently working on the extraction of cross-sections to compare to data if available and to theoretical models.





# Collaboration

## **Cyclotron Institute, Texas A&M University, USA**

Dr. Aldo Bonasera  
Dr. Marcia R. D. Rodrigues  
Dr. Victor E. Iacob  
Dr. Ninel Nica  
Dr. Brian Roeder  
Dr. Gabriel C. Tabacaru  
K. Wang  
J. Romo

## **University of Athens, Greece**

Dr. George Souliotis

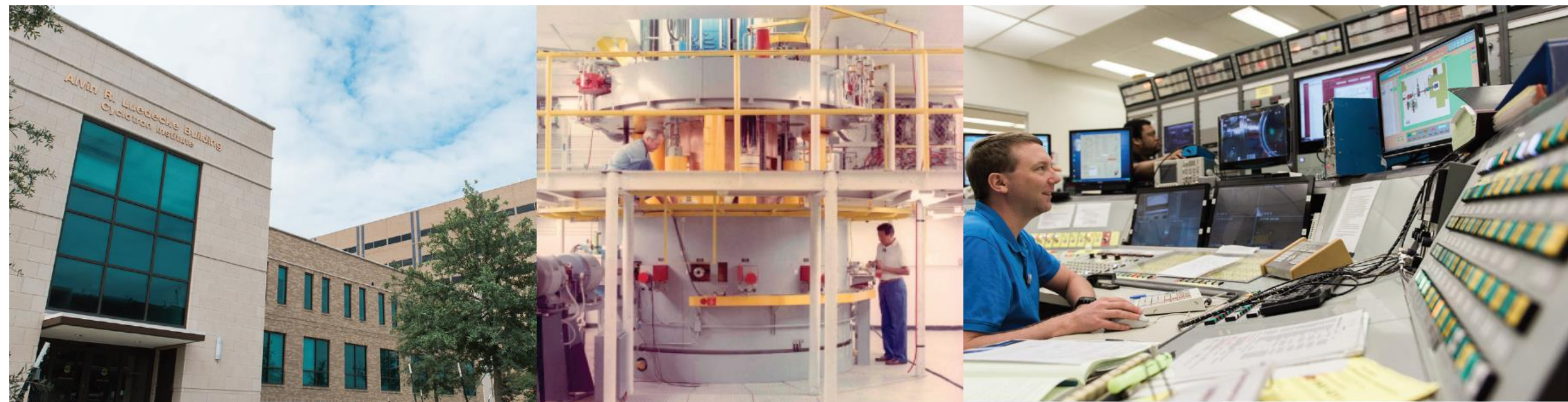
## **HMRI, Houston Methodist Hospital, USA**

Dr. Max Yu

Dr. P. Zanotti-Fregonara

## **Praire View, Texas A&M University, USA**

Dr. Justin Mabilia



Thank you for your attention

# Back up slides

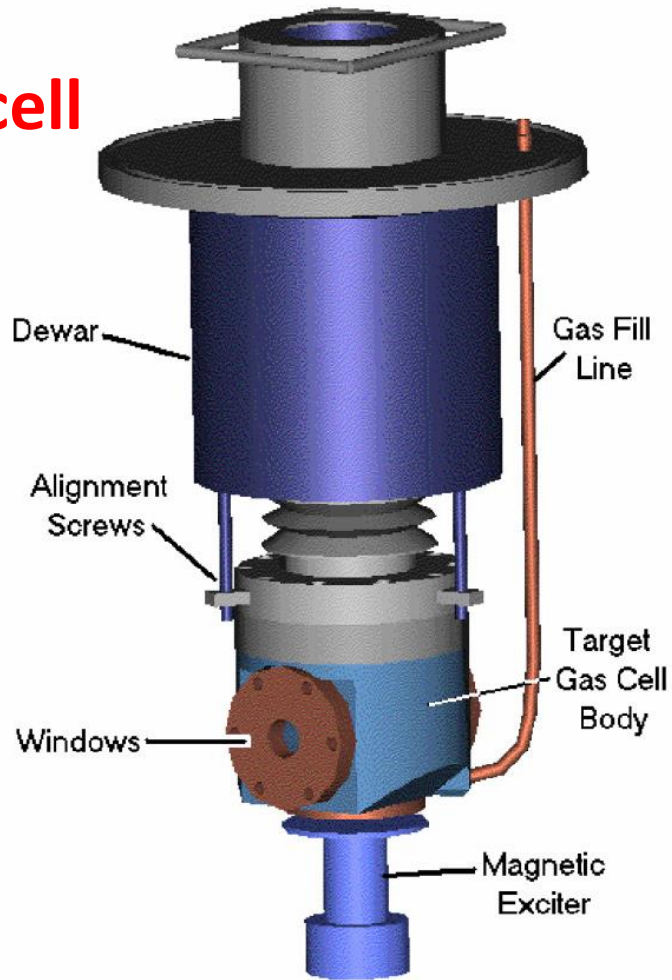


Reactors	Fission of $^{235}\text{U}$	$n + ^{235}\text{U} \rightarrow ^{99}\text{Mo} + xn + \text{other fission products}$
	Neutron activation of $^{98}\text{Mo}$	$n + ^{98}\text{Mo} \rightarrow ^{99}\text{Mo}$
Accelerators	Photo-fission of $^{238}\text{U}$	$\text{Photon} + ^{238}\text{U} \rightarrow ^{99}\text{Mo} + xn + \text{other fission products}$
	$^{100}\text{Mo}$ transmutation	$\text{Photon} + ^{100}\text{Mo} \rightarrow ^{99}\text{Mo} + n$
	Direct $^{99\text{m}}\text{Tc}$ production	$\text{P} + ^{100}\text{Mo} \rightarrow ^{99\text{m}}\text{Tc} + 2n$

Table 3. The various technological options for the production of  $^{99\text{m}}\text{Tc}/^{99}\text{Mo}$

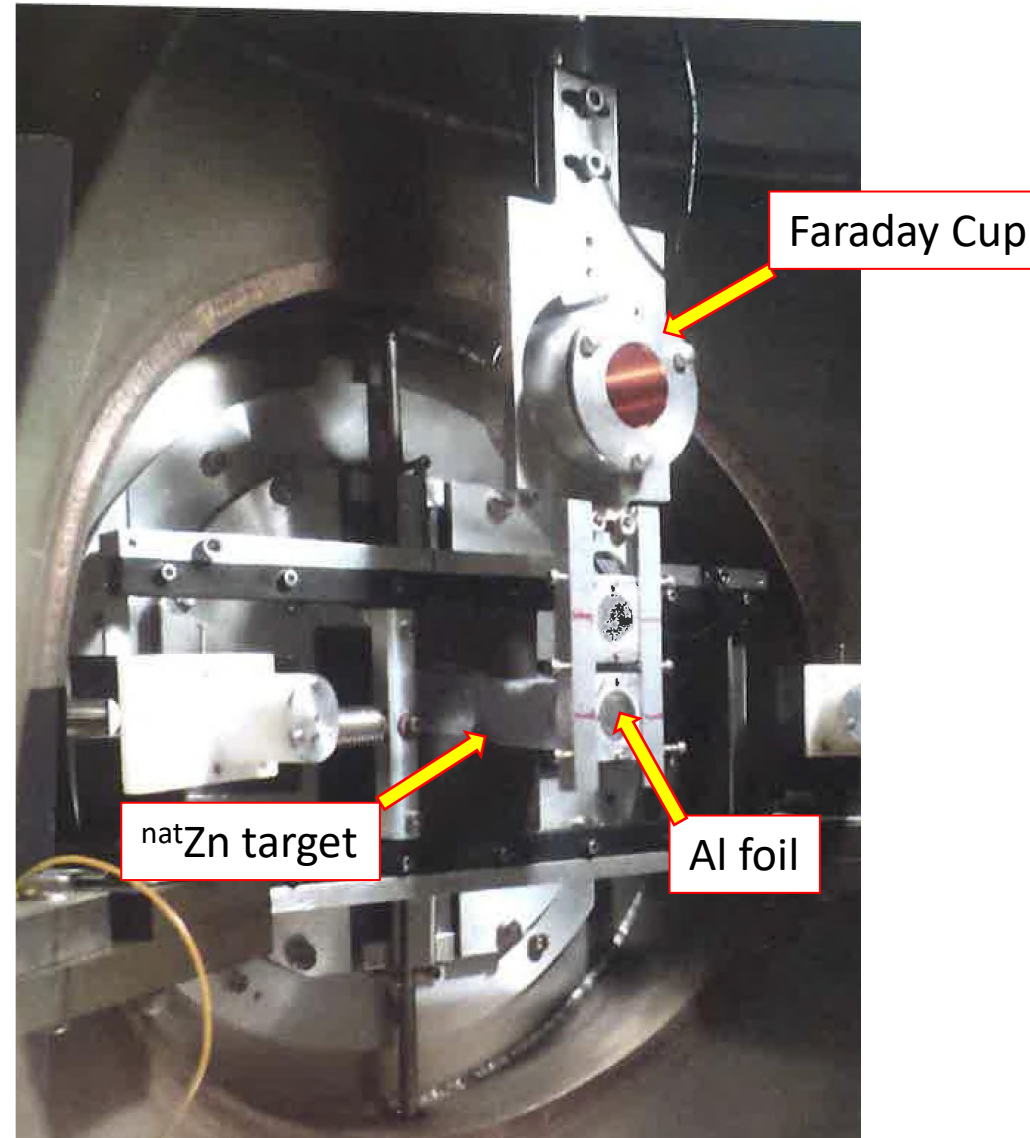
# Experimental Set up

## Cryogenic gas cell

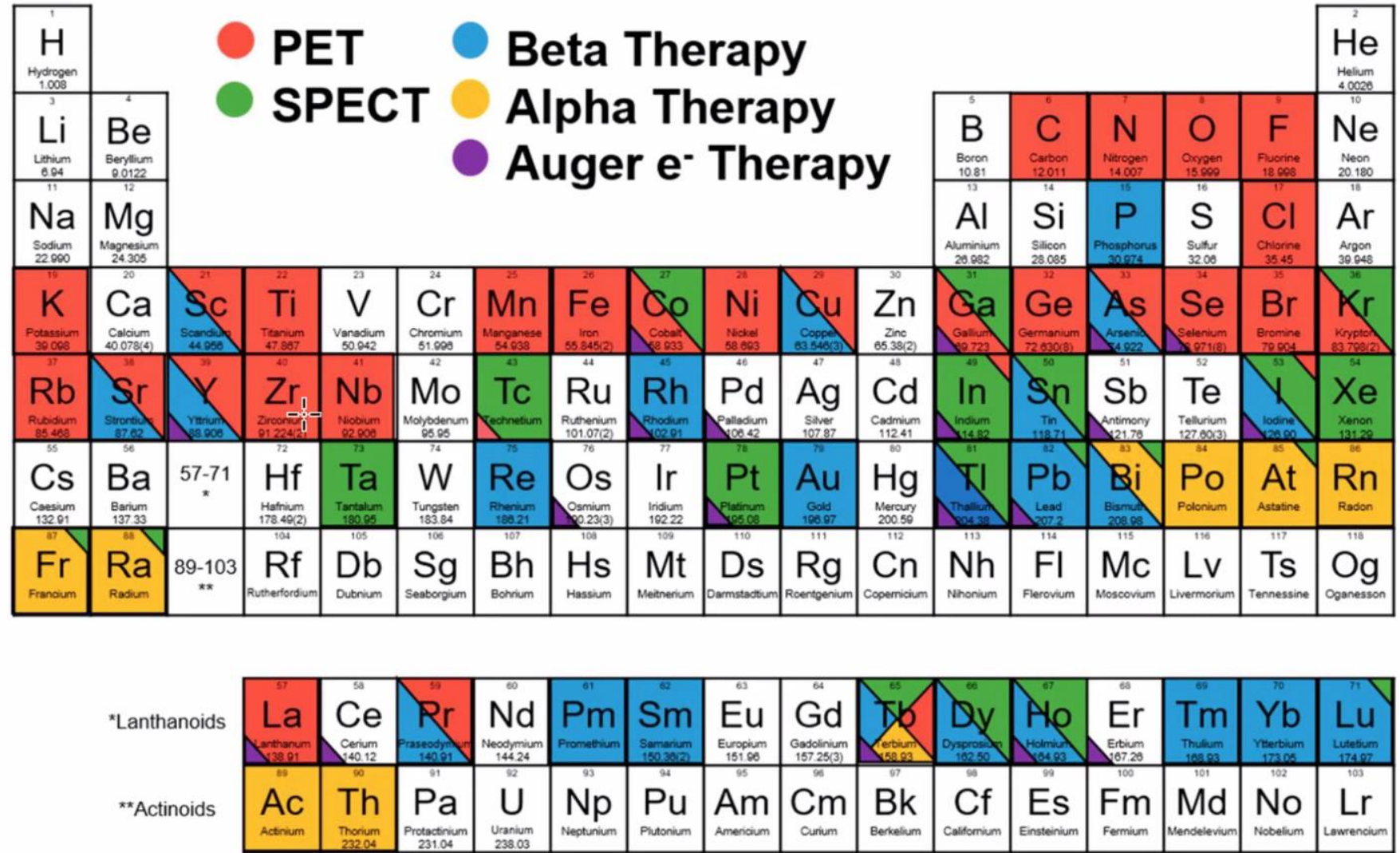


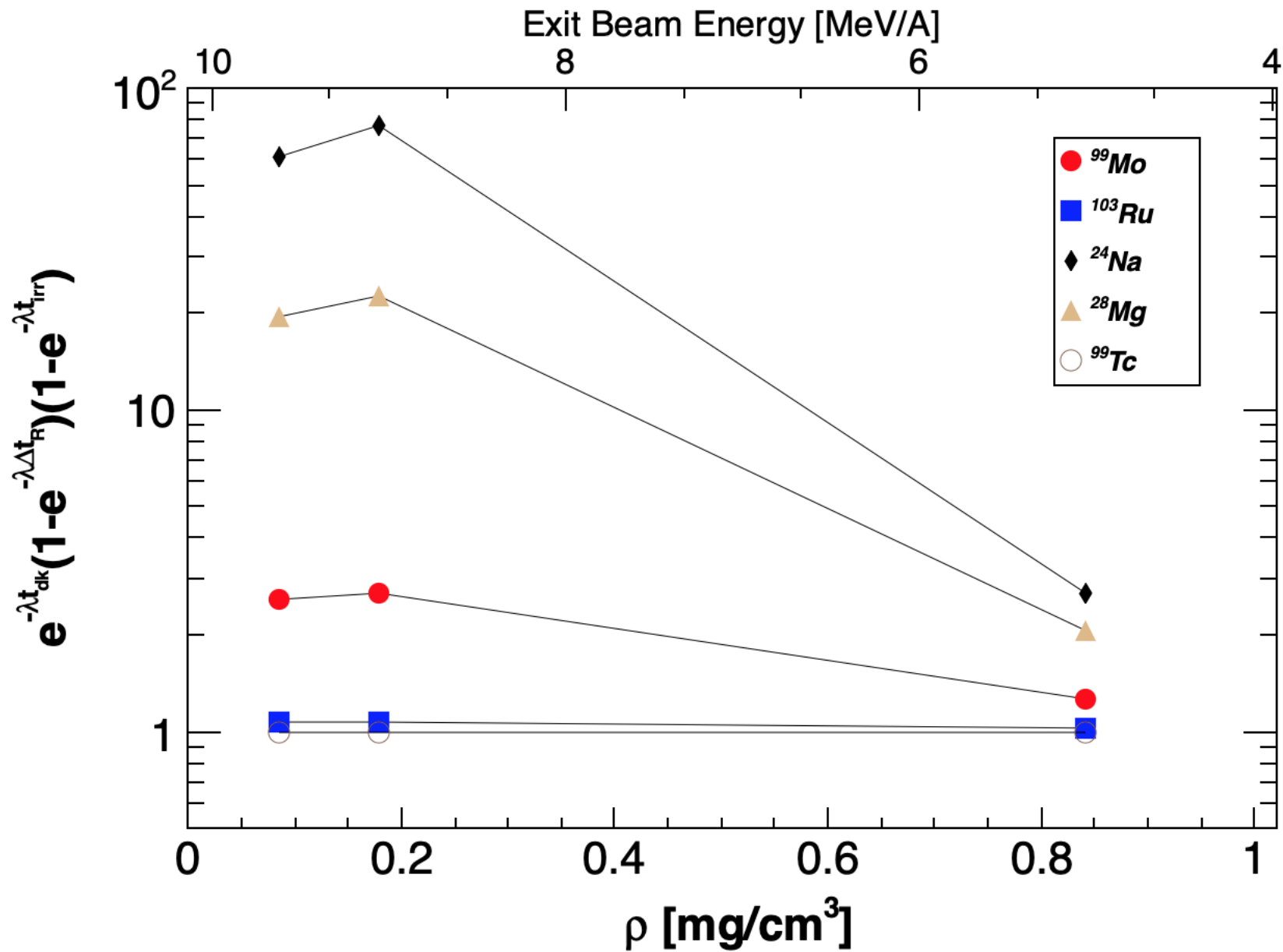


# Experimental Set up



# Medical Isotopes: Metals offer many options









## Energy threshold and Q-value

For an interaction of projectile nuclei of mass  $m_1$ , and target nuclei of mass  $m_2$ , the Q-value ( $Q$  [MeV]) is defined as:

$$Q = m_1 \cdot c^2 + m_2 \cdot c^2 - \sum_i m_i \cdot c^2 \quad (2.13)$$

where  $m_i$  denotes the mass of  $i$ -th produced nucleus. Obviously, for elastic scattering,  $Q = 0$ . For other processes, two cases are possible:

- $Q < 0$  then the reaction is endo-energetic, and the final system has higher mass (that has to be delivered in the form of kinetic energy of colliding particles),
- $Q > 0$  then the reaction is exo-energetic, and the final system has lower mass.

The exo-energetic reaction may occur at any energy. For the endo-energetic reaction, a laboratory energy limit, called threshold energy ( $E_{thr}$ ), sets the minimum energy needed for the reaction to occur. The laboratory energy threshold can be calculated from the conservation of momentum and energy to link it with Q-value:

$$E_{thr} = |Q| \cdot \left( 1 + \frac{m_1}{m_2} \right) \quad (2.14)$$