

Fission in high-energy proton-induced spallation reactions - A progress report*

N.G. Nicolis¹, A.Asimakopoulou² and G.A. Souliotis²,

¹Department of Physics, The University of Ioannina, Ioannina
45110, Greece

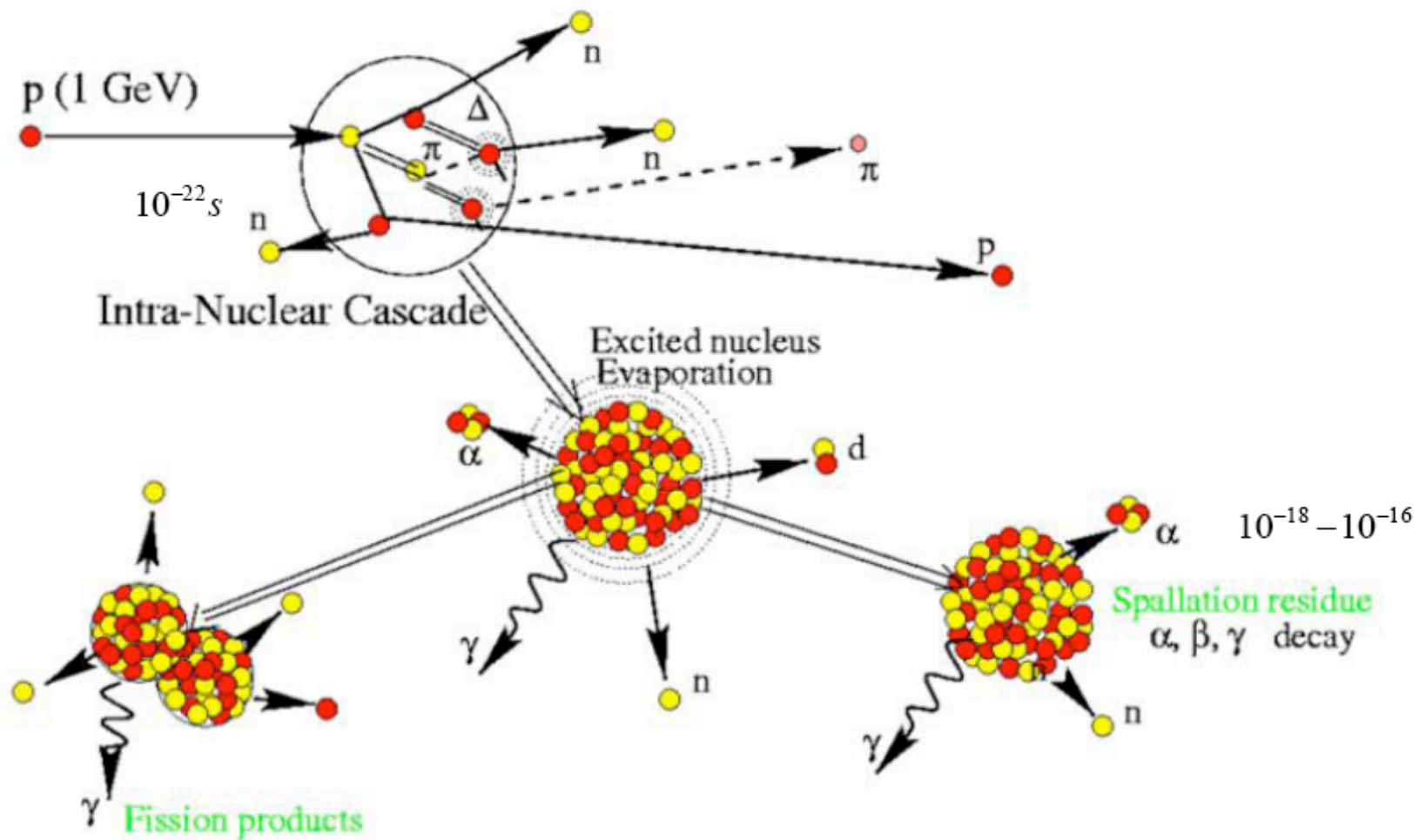
²Laboratory of Physical Chemistry, Department of Chemistry,
University of Athens, Athens, Greece

* Presented at the 4th Workshop of the Hellenic Institute of Nuclear Physics, University of Ioannina, May 5-6, 2017.

Motivation for studying spallation reactions

- Accelerator-driven systems (ADS)
- Transmutation of nuclear waste
- Spallation neutron sources
- Production of rare isotopes
- Interaction of cosmic rays with interstellar bodies (astrophysics)
- Radiation damage in space
- Testing ground for high-energy nuclear reaction models (above 150-200 MeV)

The two phases of a proton-induced spallation reaction

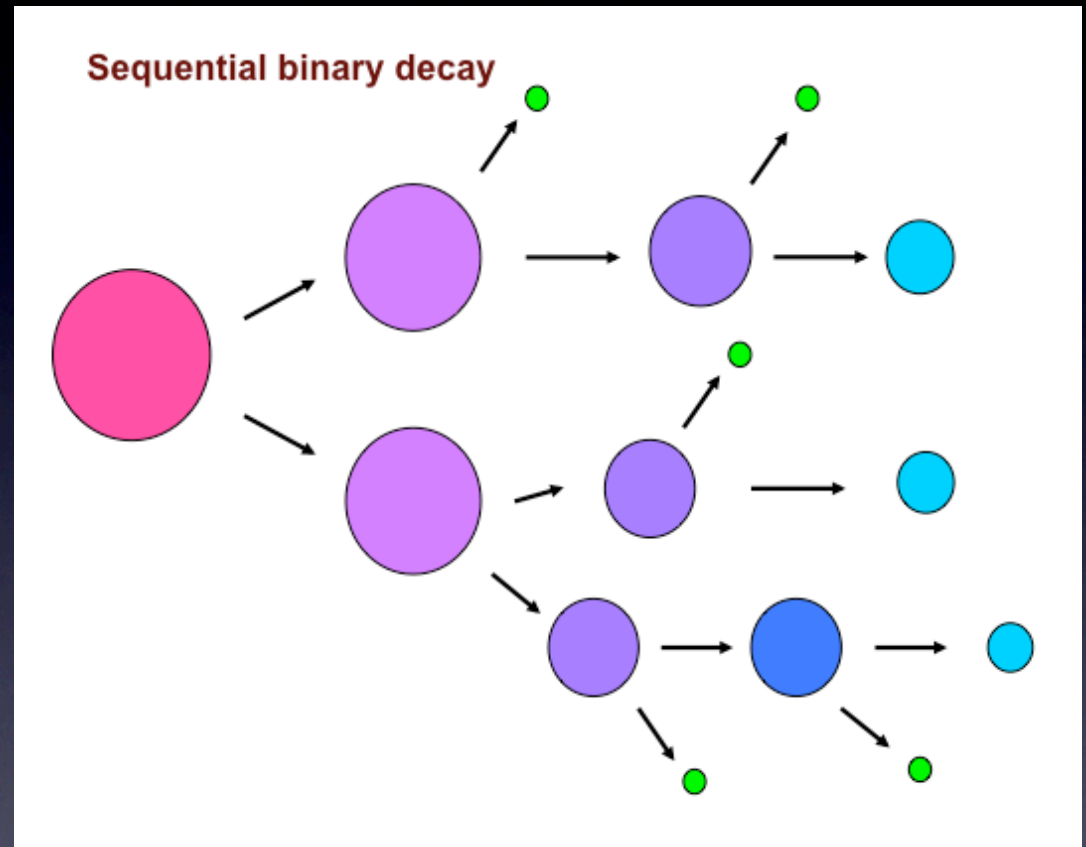
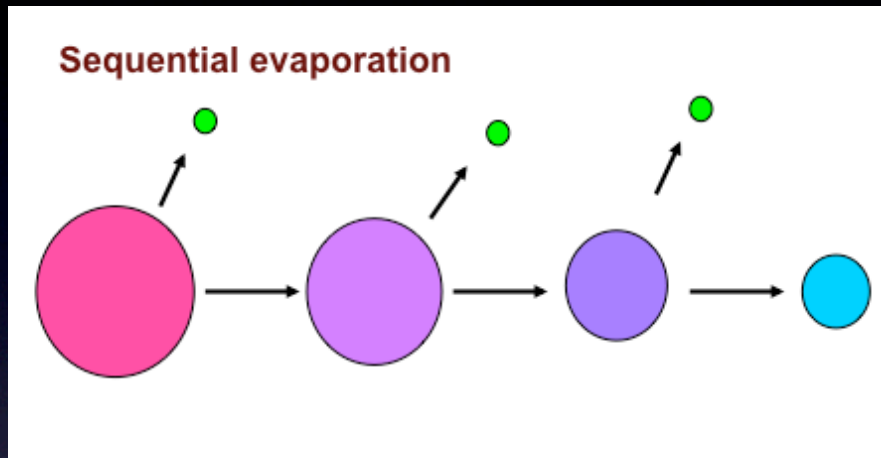


The INC phase (ISABEL code)

- Continuous medium
- Diffuse nuclear surface
- Linear trajectory between collisions
- Free N-N cross sections
- Allows for inelastic N-N collisions
- Collision criterion based on the mean-free-path
- Full Pauli blocking: Interactions resulting in nucleon falling below the Fermi sea are forbidden

Ref: Y.Yariv and Z. Fraenkel, Phys. Rev. C20, 2227 (1979); Phys. Rev. C24, 488 (1981)

The equilibrium decay phase



SMM (Botvina 2008)

MECO (Nicolis 2008)

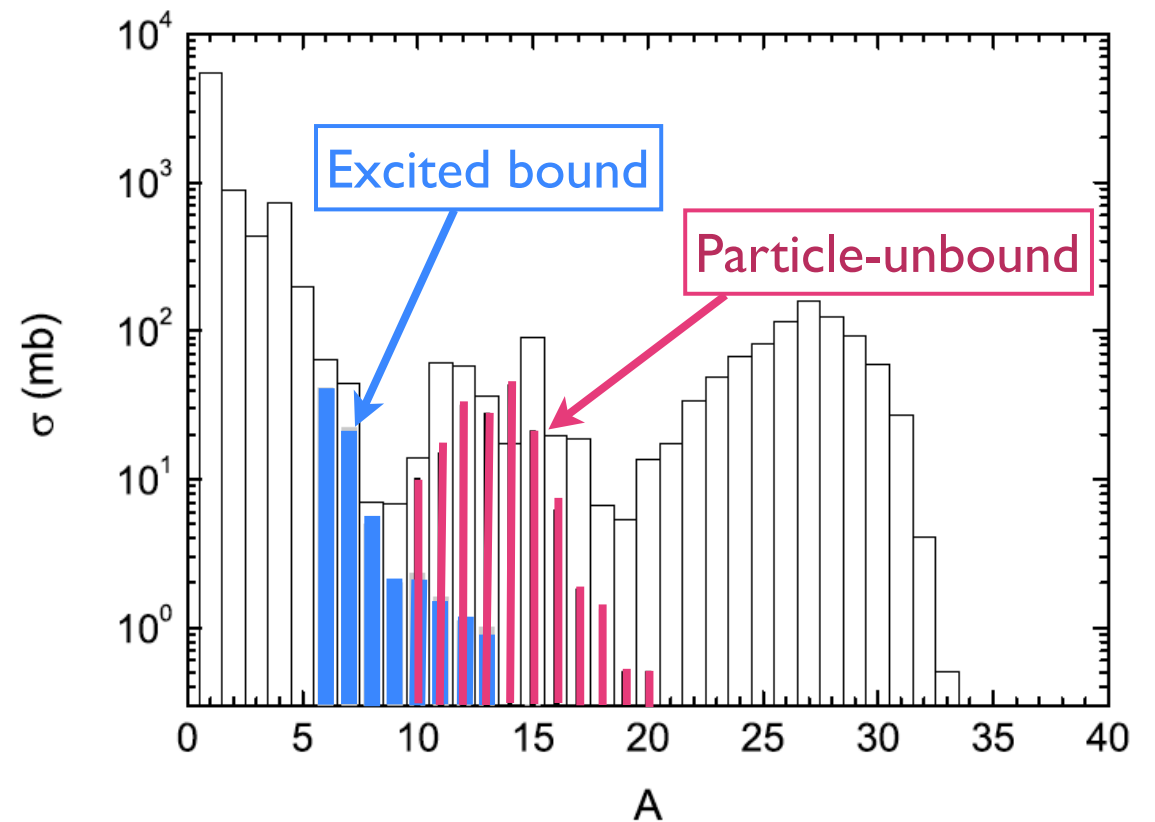
The code **MECO** (**M**ulti-sequential **E**vaporation **C**ode)

- Monte-Carlo code
- The equilibrium decay of excited nuclei is described as a sequence of binary processes involving emission of fragments in their ground, excited bound and unbound states.
- Any number of user-defined channels can be used.
- Emission of nucleons up to symmetric mass divisions can be used in a generalized Weisskopf evaporation formalism.
- Complete description of reaction systems with an effective fissility below the Businaro-Galone point.

MECO run with 179 decay channels

$^{40}\text{Ar}^*$ (180 MeV)

γ (E1)	
n	g.s.
$1\text{-}^3\text{H}$	g.s.
$3,5\text{He}$	g.s.
$6\text{-}^8\text{Li}$	g.s. + 7 e.s.
$7\text{-}^{10}\text{Be}$	g.s. + 7 e.s.
$8\text{-}^{12}\text{B}$	g.s. + 18 e.s. + c.s.
$11\text{-}^{16}\text{C}$	g.s. + 21 e.s. + c.s.
$13\text{-}^{18}\text{N}$	g.s. + 15 e.s. + c.s.
$15\text{-}^{21}\text{O}$	g.s. + 16 e.s. + c.s.
$16\text{-}^{20}\text{F}$	g.s. + 27 e.s. + c.s.



Ref: N.G.Nicolis, Int. Jour. Mod. Phys. E 17 (2008) 1541-1556.

Fission probability and decay width

Fission probability:

$$P_f = \frac{\Gamma_f}{\sum_j \Gamma_j + \Gamma_f}$$

Decay width for the emission of a particle j ($j=n, p, \alpha, \dots$):
(Weisskopf)

$$\Gamma_j = \frac{(2s_j + 1)m_j}{\pi^2 \rho_c(U_c)} \int_{V_j}^{U_j - B_j} \sigma_{inv}^j(E) \rho_j(U_j - B_j - E) E dE$$

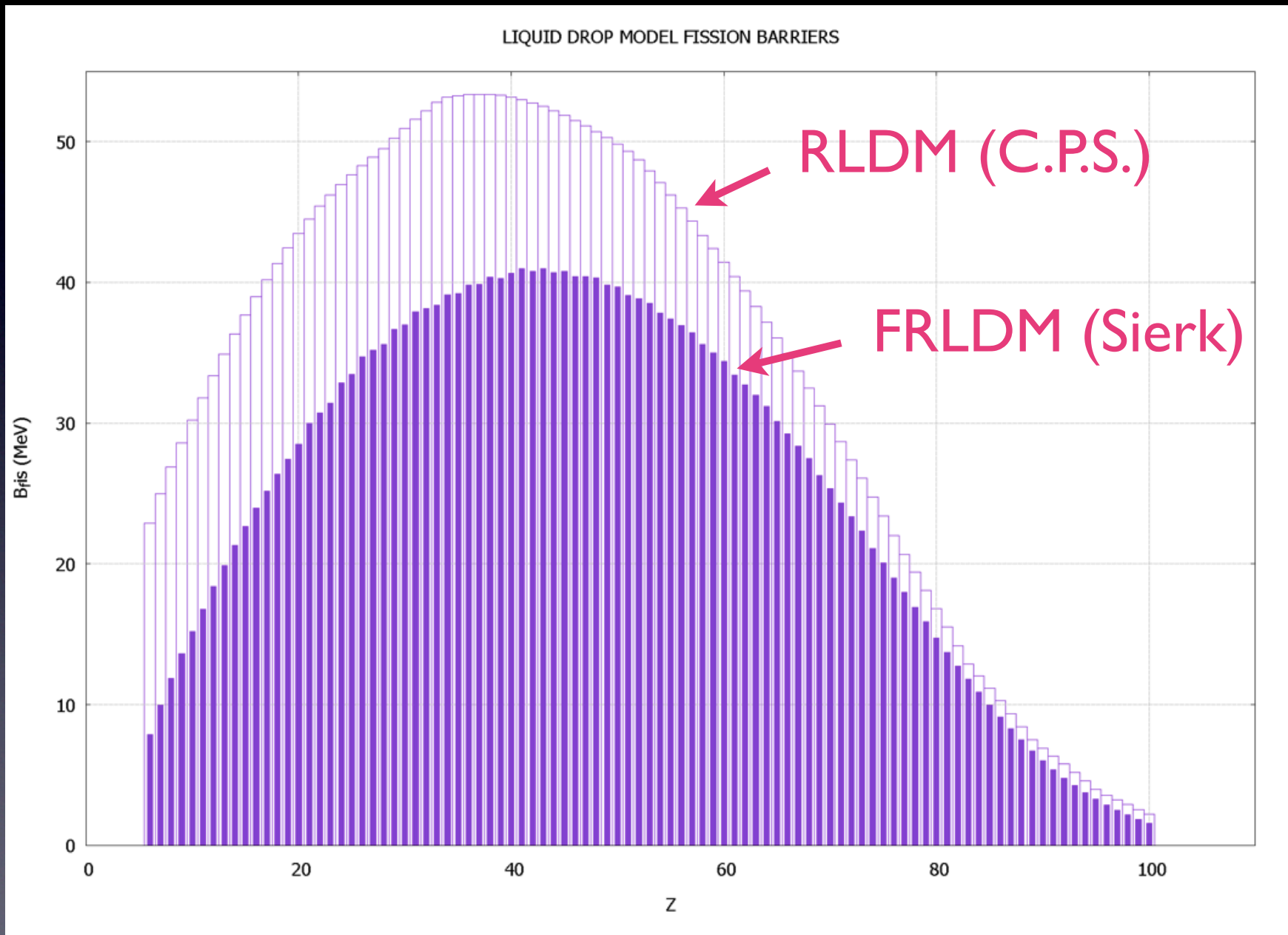
Decay width for fission (transition state theory):

$$\Gamma_f = \frac{N_{sad}}{2\pi\rho_c(U_c)}$$

or

$$\Gamma_f = \frac{1}{2\pi\rho_c(U_c)} \int_0^{U_f - B_f} \rho_f(U_f - B_f - E) dE$$

Fission barriers



Fission delay

The standard theory of fission underestimates the measured pre-fission neutron multiplicities. A fission delay is needed to account for the slowing effects of nuclear dissipation.

$$\Gamma_f = \Gamma_f^{BW} \left(\sqrt{1 + \gamma^2} - \gamma \right)$$

$$\left(\sqrt{1 + \gamma^2} - \gamma \right)$$

is the Kramers reduction factor ($\gamma \sim 5$)

$$\gamma = \frac{\beta}{2\omega_{sp}}$$

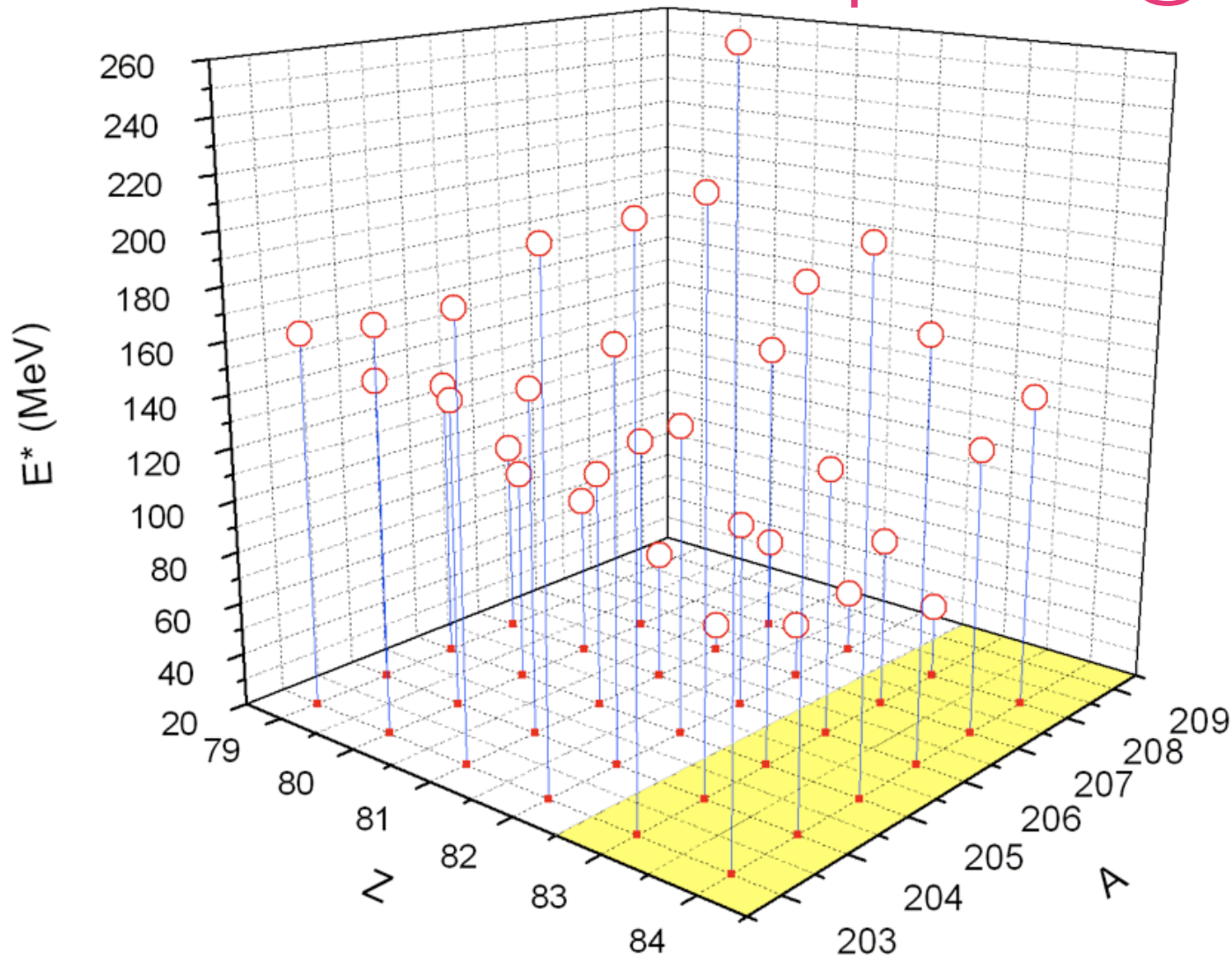
β is the reduced nuclear dissipation coefficient
 ω_{sp} is the curvature of the PES at the SP

First-order description of fission decay in MECO

- Calculate fission probability.
- If fission occurs, assume a Gaussian profile for the symmetric division with excitation energy dependent variance (G.D.Adeev 1993).
- From the conservation of energy and linear momentum we get the fragment kinetic energies.
- For each fragment, the excitation energy is divided in proportion to the fragment mass.
- Then, the particle decay of each fragment is followed.

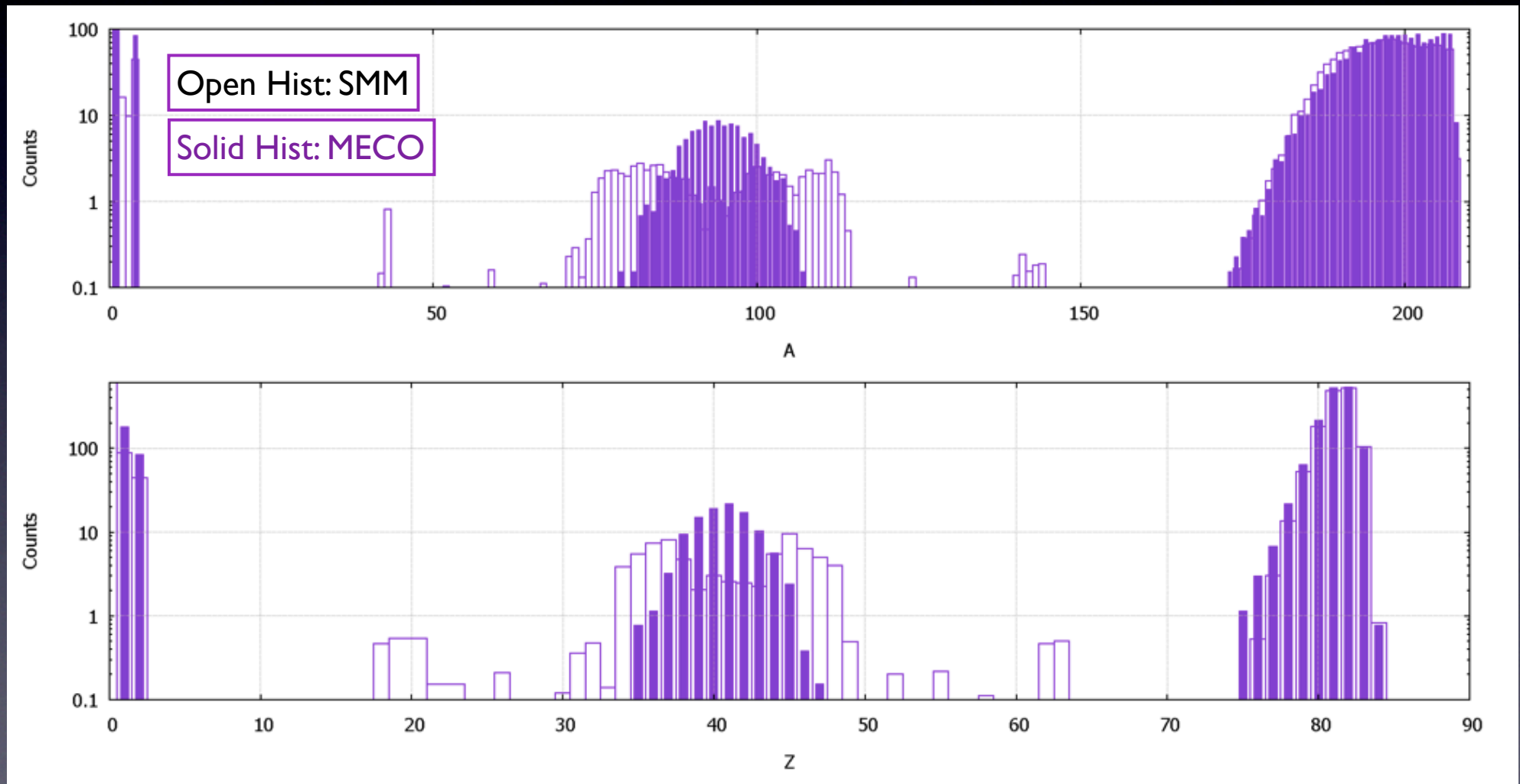
Distribution in A and Z of the average excitation energy predicted by INC

$p + {}^{208}\text{Pb}$ @ 1000 MeV

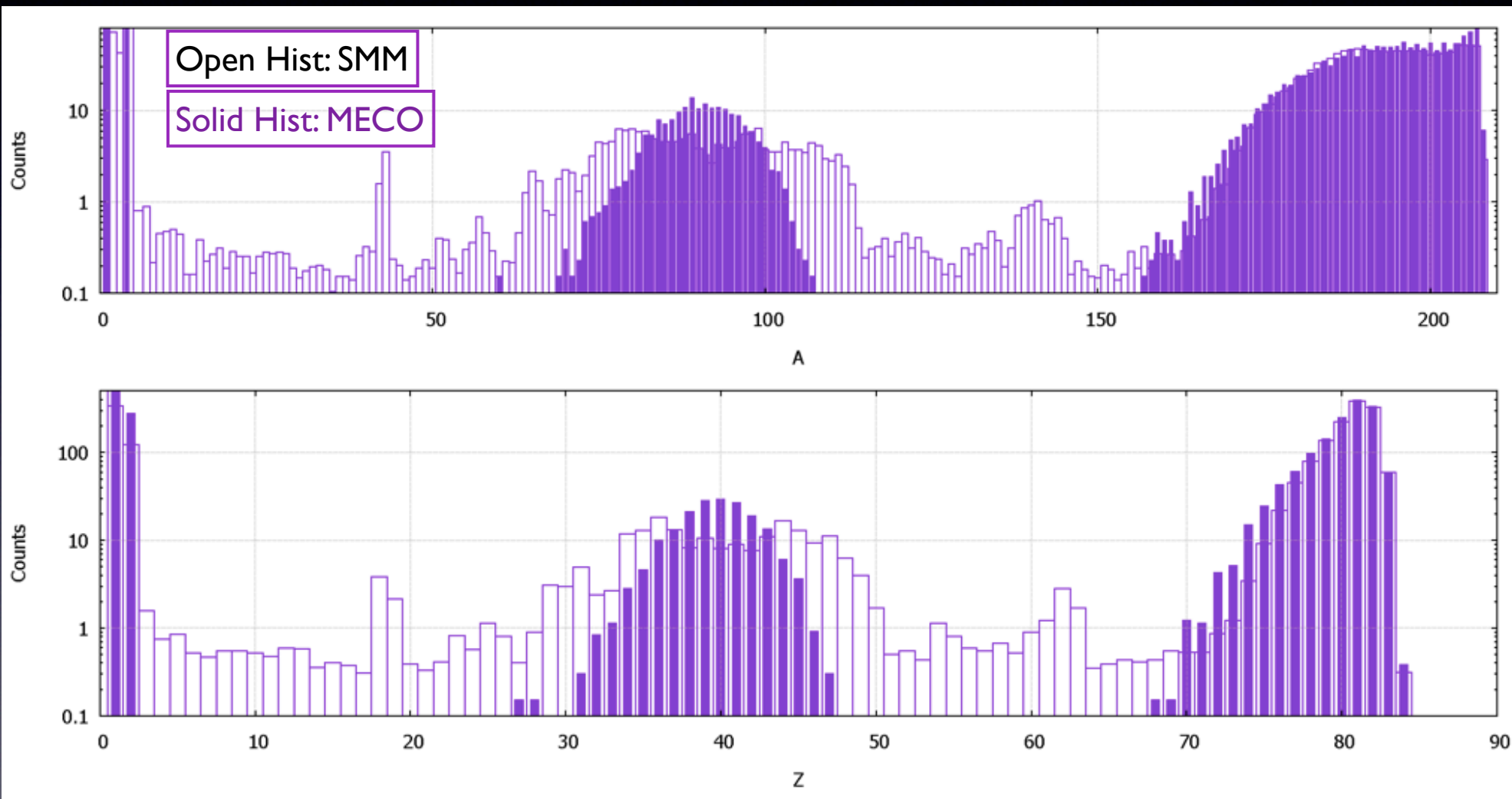


$p + {}^{208}\text{Pb}$ @ 500MeV

Calculated A and Z distributions with MECO and SMM

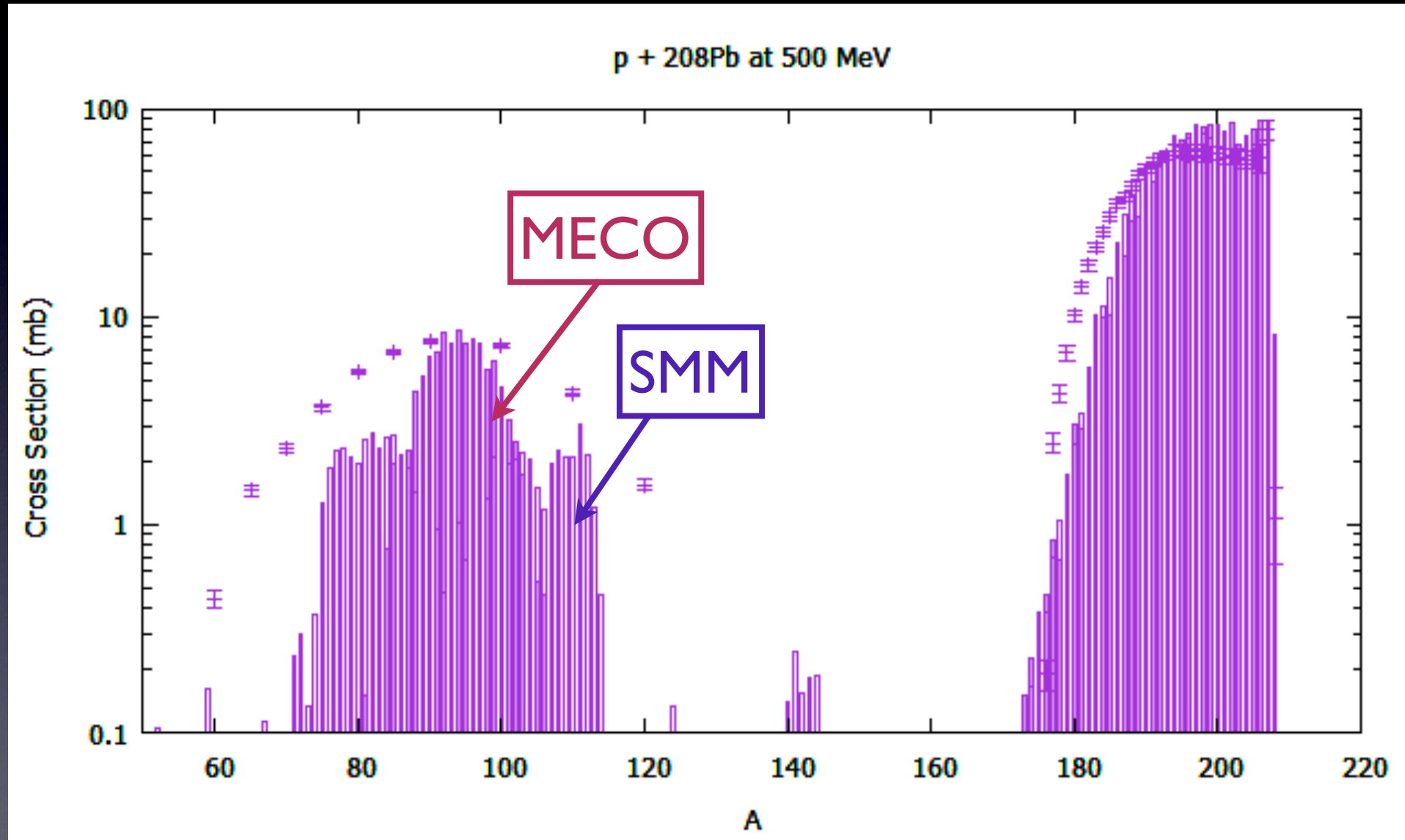


$p + {}^{208}\text{Pb}$ @ 1000MeV
Calculated A and Z distributions with MECO and SMM



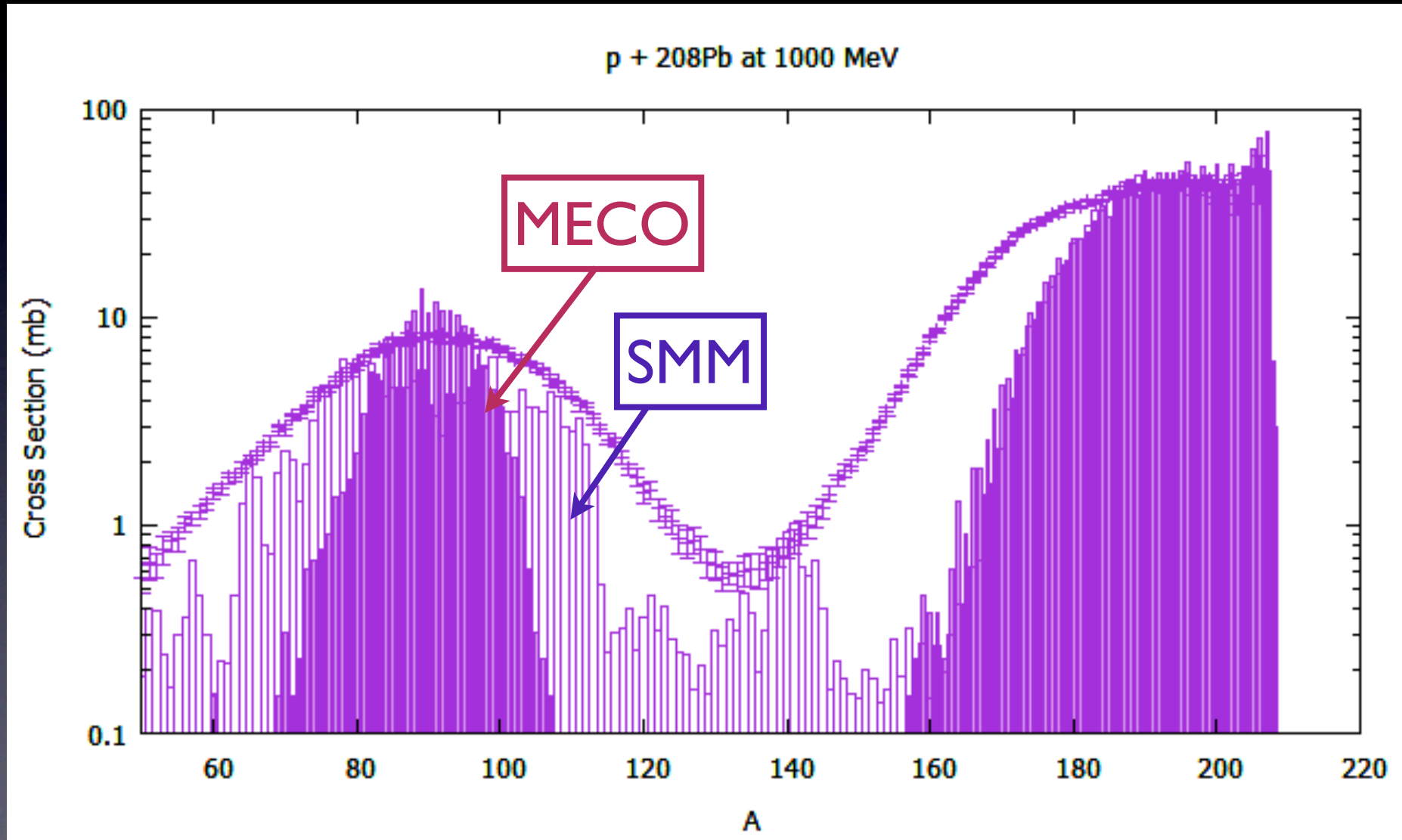
$p+^{208}\text{Pb}$ @ 500 MeV

Experimental and calculated mass distributions



$p+^{208}\text{Pb}$ @ 1000 MeV

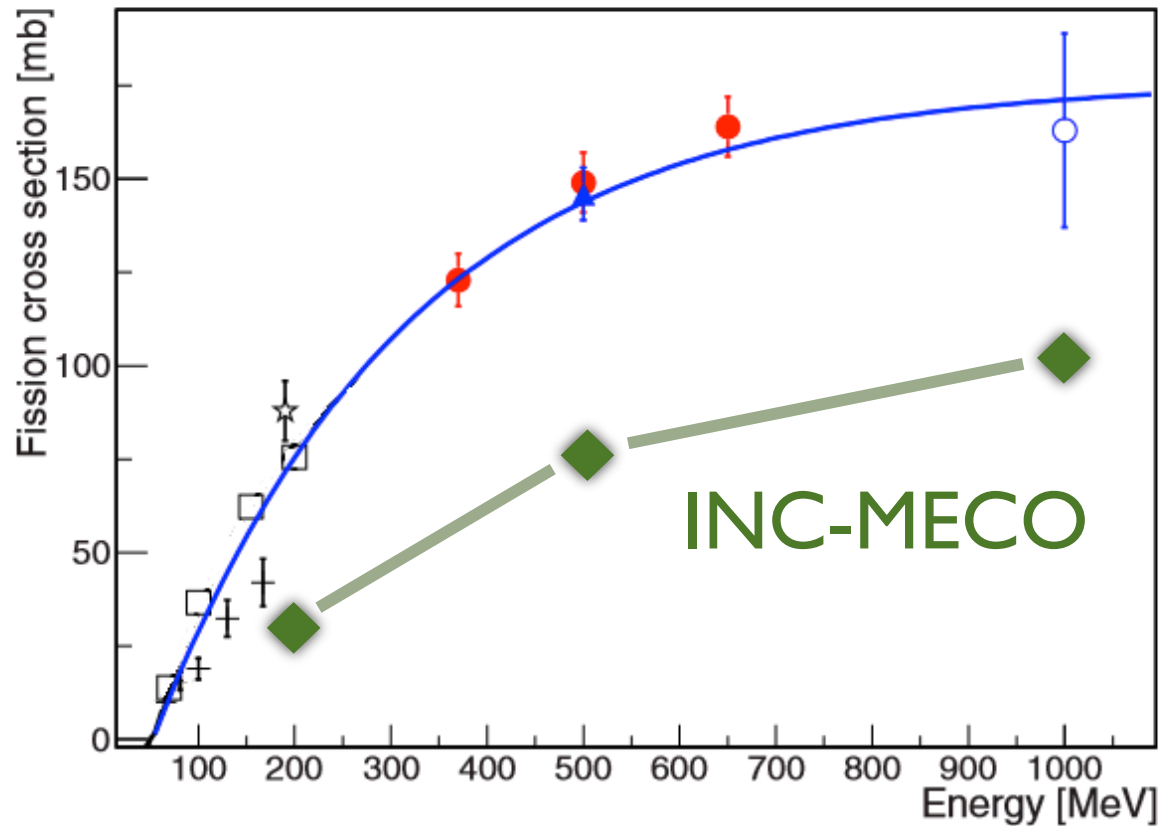
Experimental and calculated mass distributions



Description of isotopic distributions of evaporation residues

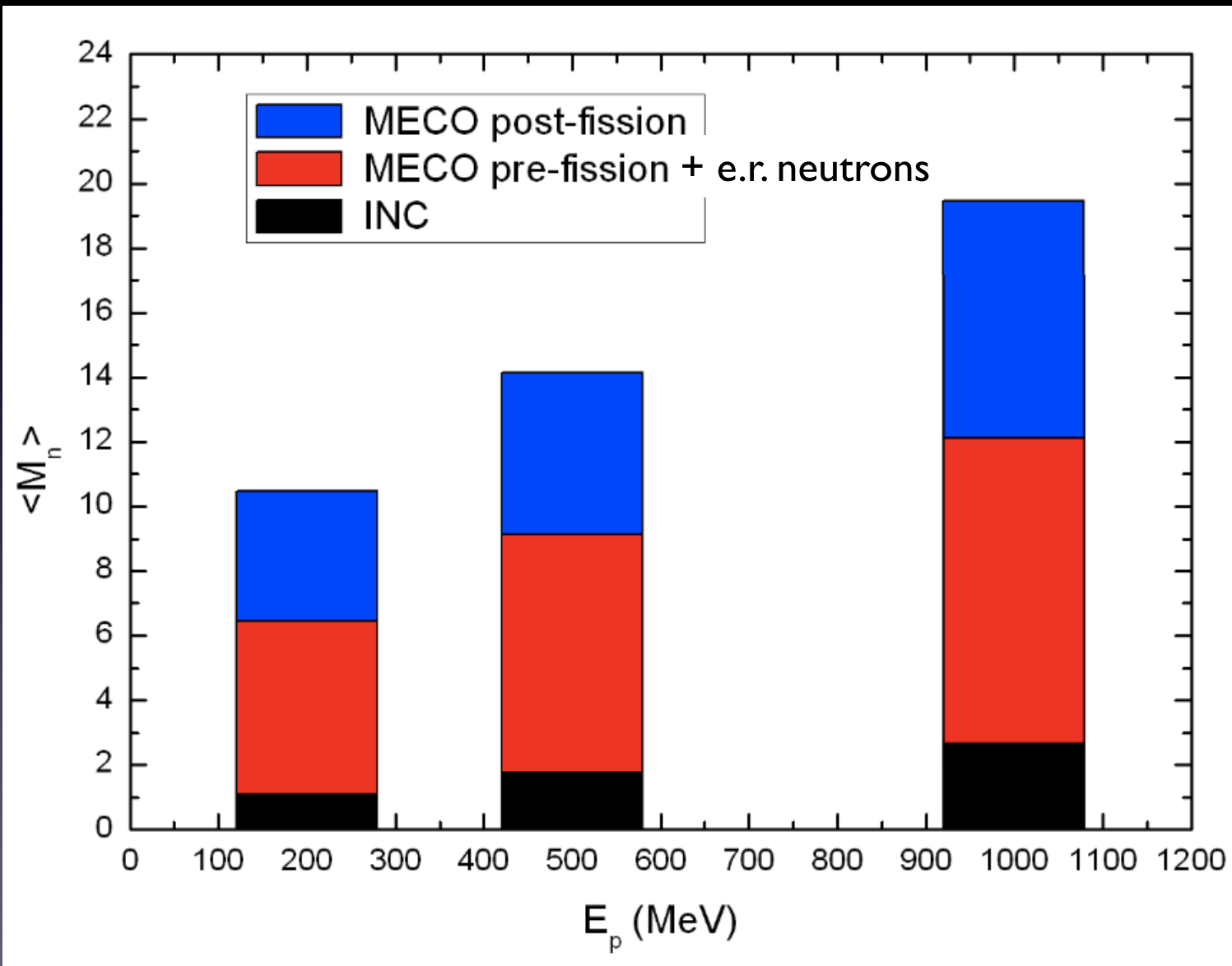
- Agreement between MECO and SMM, in most cases.
- Agreement between experimental data and calculations.
- Examples will be given this afternoon in the talk by Mrs. Aggeliki Asimakopoulou

Fission cross sections for $p + {}^{208}\text{Pb}$



- Rodriguez
- ▲ Schmidt K.-H.
- Enqvist T.
- Flerov G.N.
- + Shigaev O.E.
- ☆ Duijvestijn M.C.
- Fit

Neutron multiplicities



Summary

- We have combined the intranuclear cascade code ISABEL and the sequential binary decay code MECO in an effort to study spallation reactions induced by high-energy protons.
- Preliminary calculations of $p + {}^{208}\text{Pb}$ at bombarding energies 200, 500 and 1000 MeV were compared with experimental mass distributions, fission cross sections and neutron multiplicities for these reactions.
- For the evaporation residue mass, charge (and isotopic distributions) our calculations with MECO are consistent with the predictions of the SMM code.
- The description of the experimental fission fragment mass distributions requires a careful modeling of the temperature-dependent fission barriers.
- In this direction, additional developments in MECO are in progress.